

Crystallographic Diffraction

Laue diffraction from a three-dimensional abc lattice grating

$$\begin{aligned} a(\cos \nu_1 - \cos \mu_1) &= \mathbf{a} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = h\lambda \\ b(\cos \nu_2 - \cos \mu_2) &= \mathbf{b} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = k\lambda \\ c(\cos \nu_3 - \cos \mu_3) &= \mathbf{c} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = l\lambda \end{aligned}$$

Max Laue, Walther Friedrich, and Paul Knipping (1912).

Bragg reflection from families of parallel hkl lattice planes

$$2d_{hkl} \sin \theta = n\lambda, \quad 2\left(\frac{d_{hkl}}{n}\right) \sin \theta = \lambda, \quad 2d_{nhnknl} \sin \theta = \lambda$$

William Henry and William Lawrence Bragg (1913). (Father and son)

Integrated Bragg reflection intensities

$$\rho = \frac{E\omega}{I_0} = kALp|F_{hkl}|^2 = \left(\frac{e^2}{mc^2}\right)^2 \lambda^3 \left(\frac{v_{\text{xtal}}}{V_{\text{cell}}}\right)^2 \left[\int_{v_{\text{xtal}}} e^{-\mu(t_0+t_1)} dv \right] \frac{1}{\sin 2\theta} \left(\frac{1}{2} + \frac{1}{2} \cos^2 2\theta\right) |F_{hkl}|^2$$

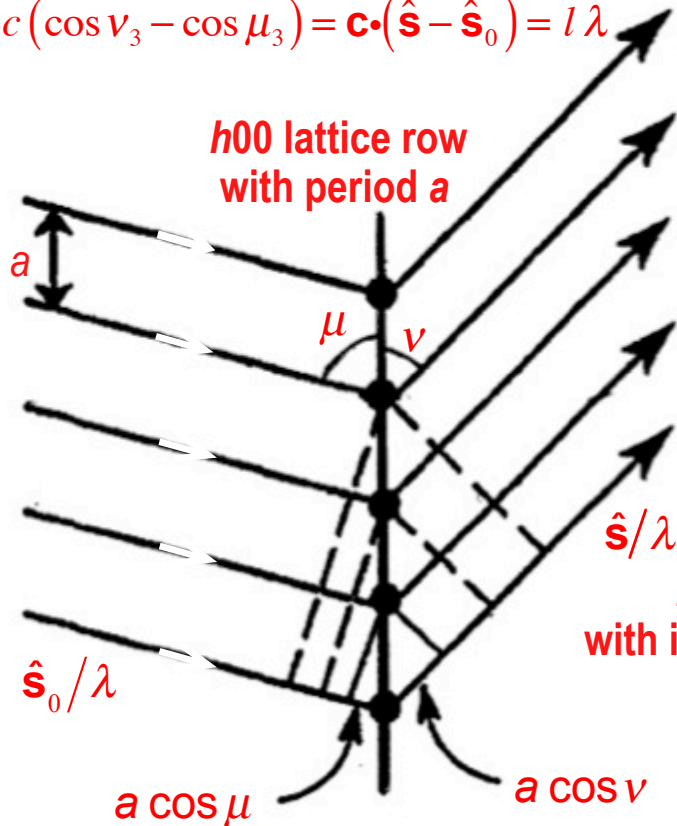
Charles G. Darwin (1914). (Grandson of the author of the theory of evolution)

Crystallographic diffraction

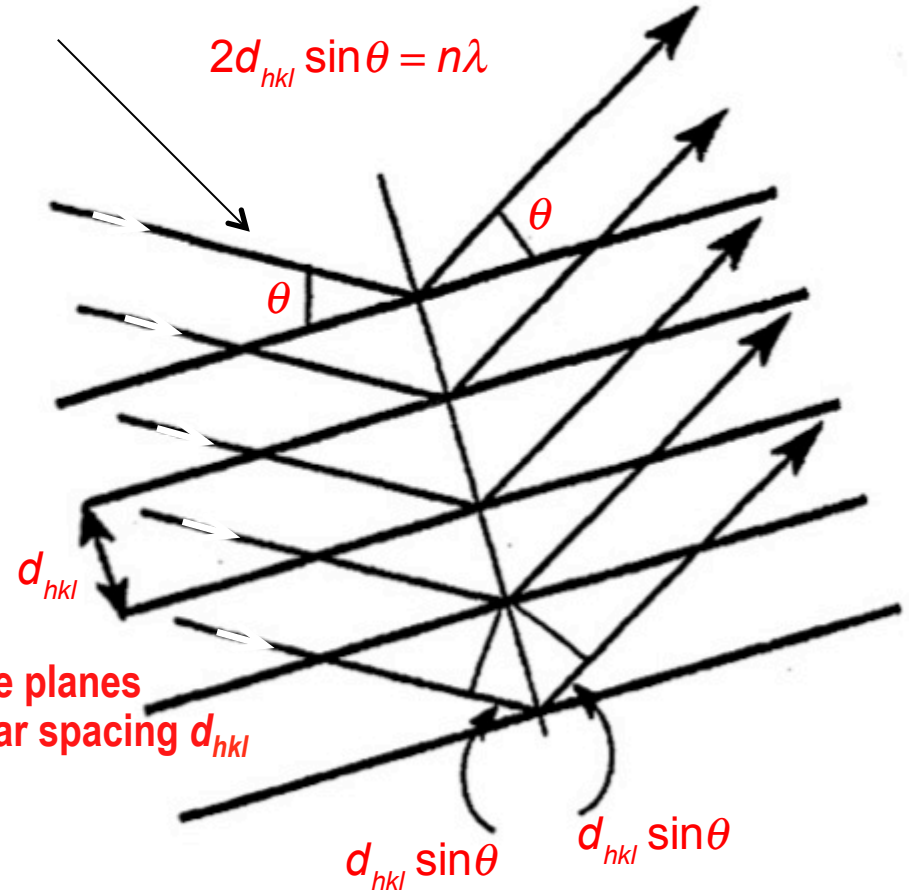
$$a(\cos \nu_1 - \cos \mu_1) = \mathbf{a} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = h\lambda$$

$$b(\cos \nu_2 - \cos \mu_2) = \mathbf{b} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = k\lambda$$

$$c(\cos \nu_3 - \cos \mu_3) = \mathbf{c} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = l\lambda$$

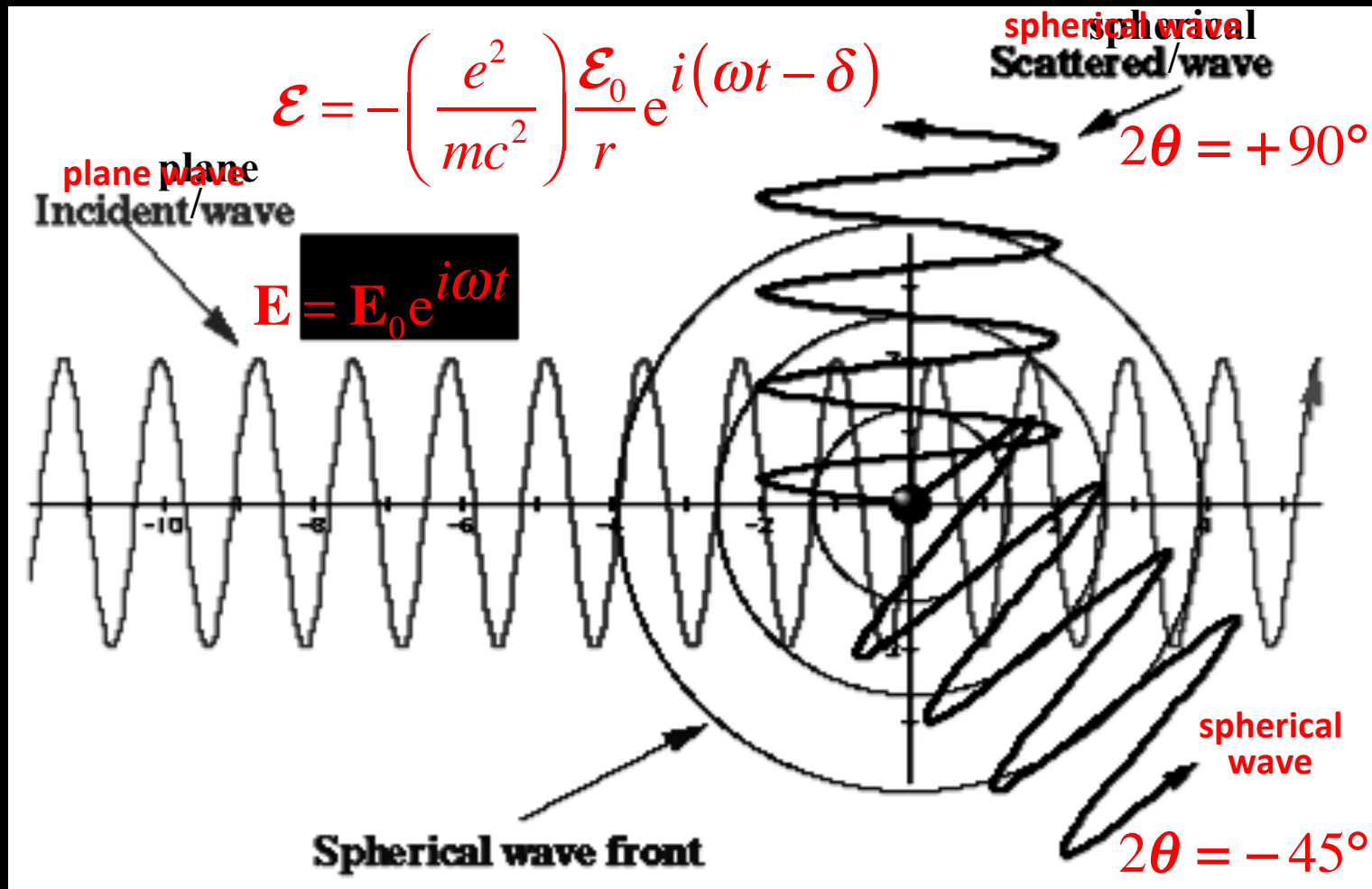


Laue transmission



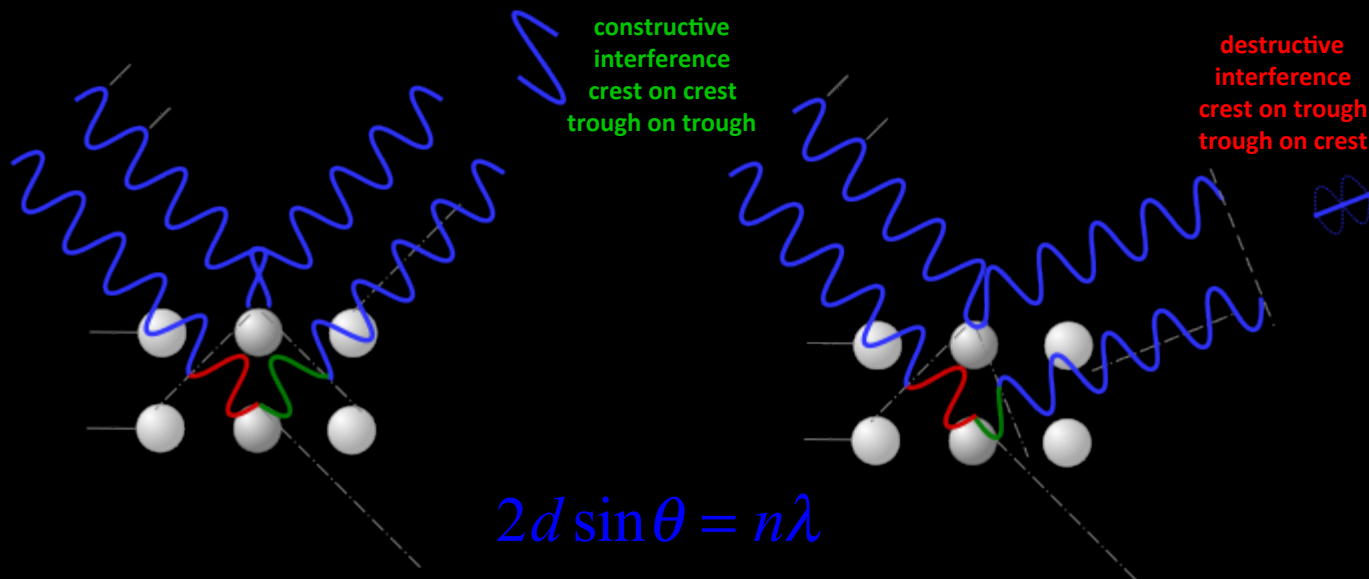
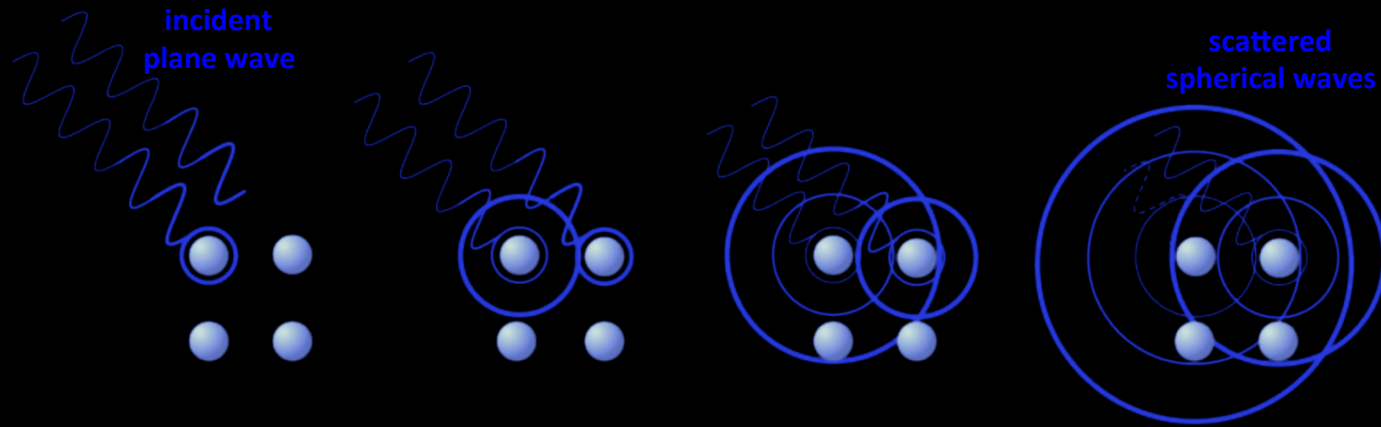
Bragg reflection

Thomson scattering of an X-ray wave due to driven oscillation of an electron

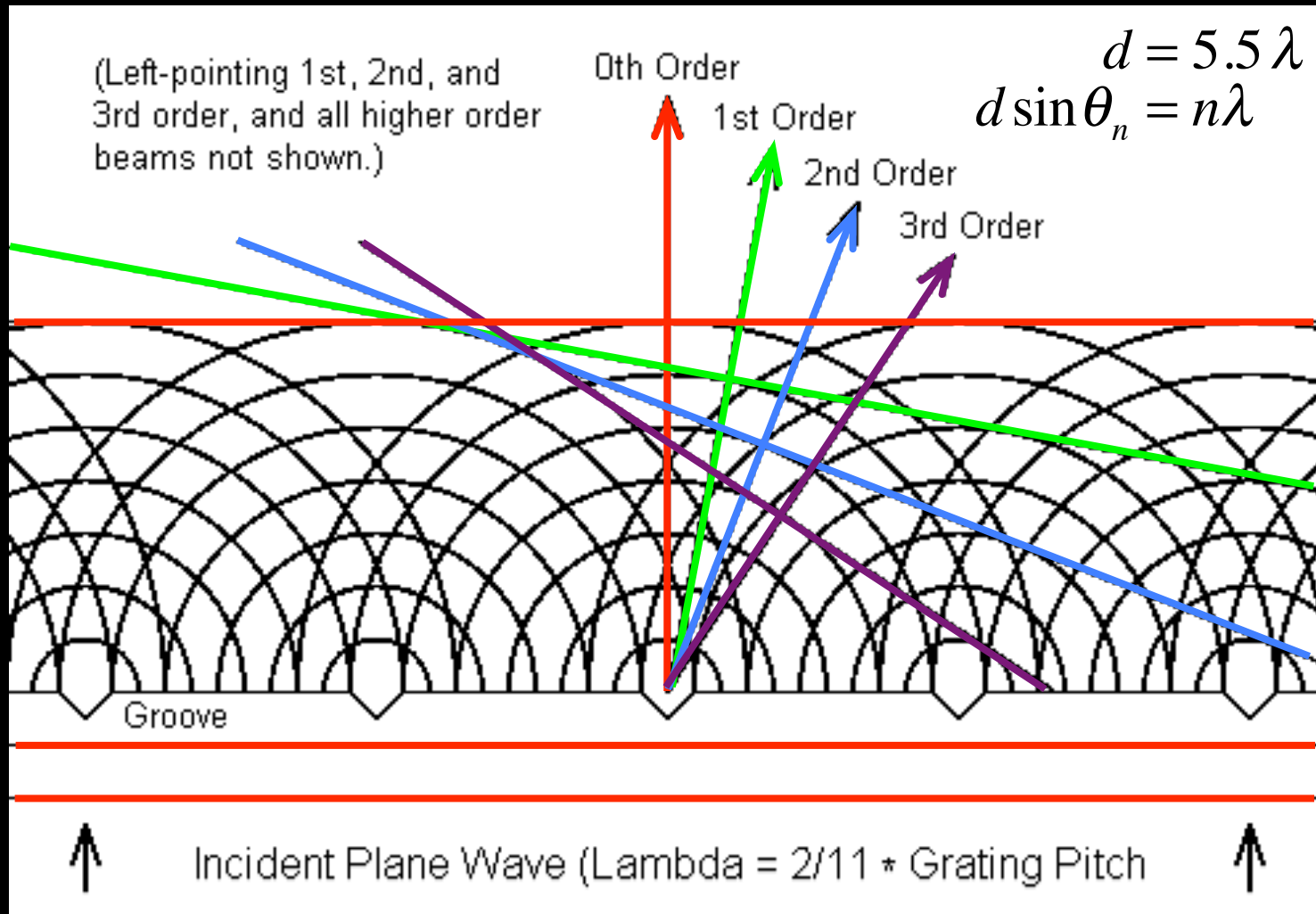


Elastic scattering of an electromagnetic wave

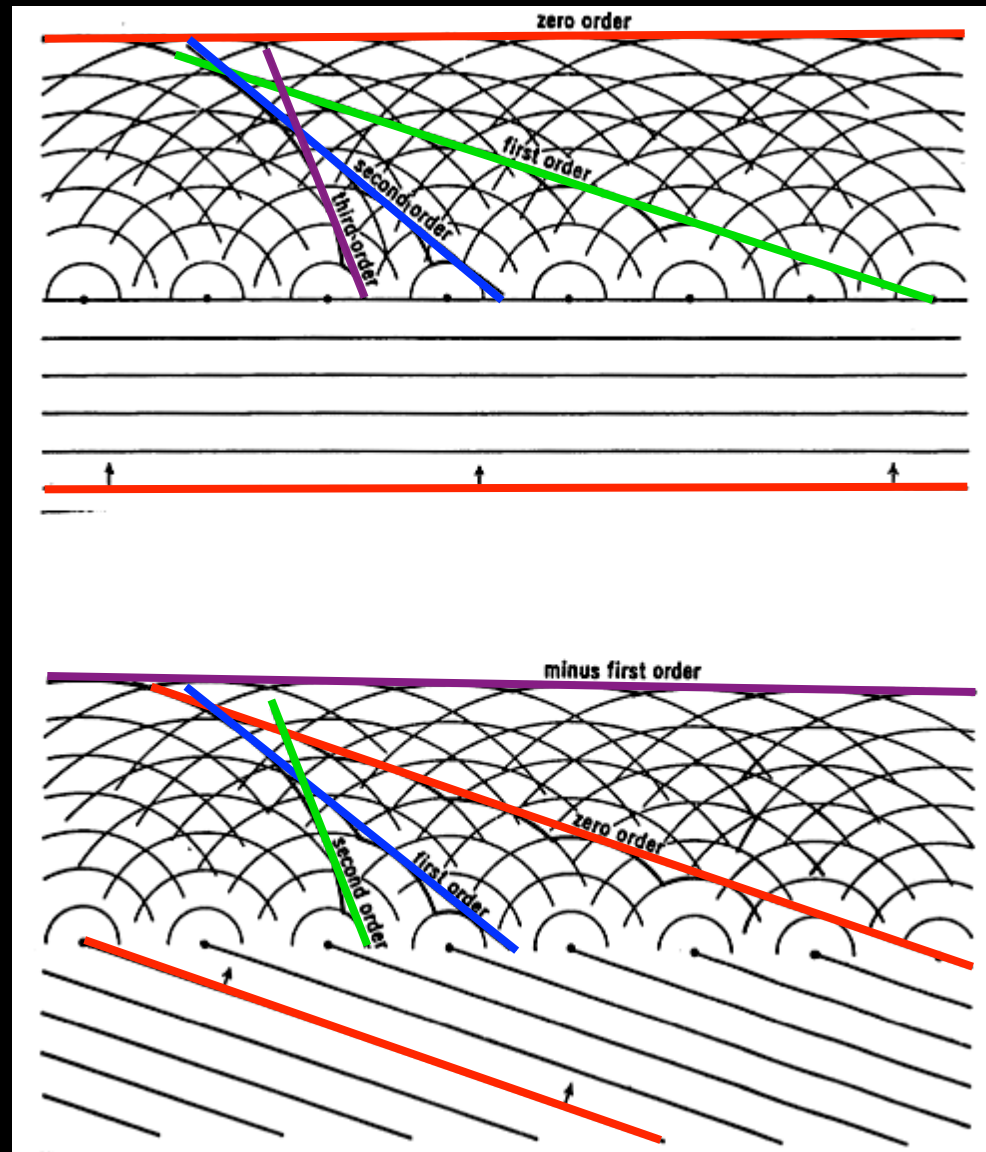
Rayleigh-Thomson scattering Laue diffraction and Bragg reflection



Huygens construction for a diffraction grating with $d = 5.5 \lambda$



Huygens constructions for diffraction by a grating or by a lattice row



\vec{a}

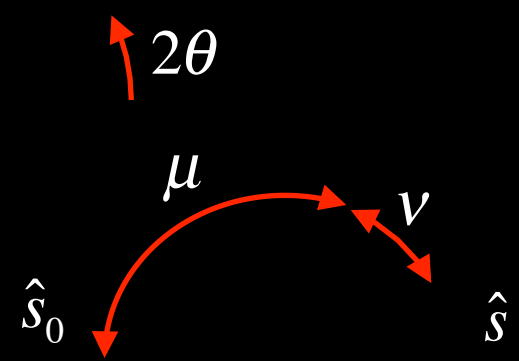
$$\vec{k}_0 + \vec{S} = \vec{k}$$

$$\vec{k} = \hat{s} / \lambda$$

$$\vec{\mathcal{E}} = \vec{\mathcal{E}}_0 e^{i(\omega t - \delta)}$$

$$\vec{E} = \vec{E}_0 e^{i\omega t}$$

$$\vec{k}_0 = \hat{s}_0 / \lambda$$

 λ \vec{a} \hat{s} \hat{s}_0 $\vec{a} \cdot \hat{s}$

$$|\hat{s}_0| = |\hat{s}| = 1$$

$$\hat{s}_0 \cdot \hat{s} = \cos 2\theta$$

$$\vec{S} = \vec{k} - \vec{k}_0 = \frac{\hat{s} - \hat{s}_0}{\lambda}$$

$$|\vec{S}| = 2 \left(\frac{\sin \theta}{\lambda} \right)$$

$$\text{path 2} = |\hat{s}_0| + \vec{a} \cdot \hat{s}_0 + |\hat{s}|$$

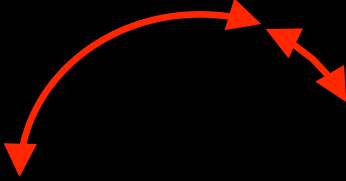
$$\text{path 1} - \text{path 2} = \vec{a} \cdot \hat{s} - \vec{a} \cdot \hat{s}_0 = \vec{a} \cdot (\hat{s} - \hat{s}_0)$$

$$\vec{a} \cdot \hat{s}_0 = |\vec{a}| \cos \mu$$

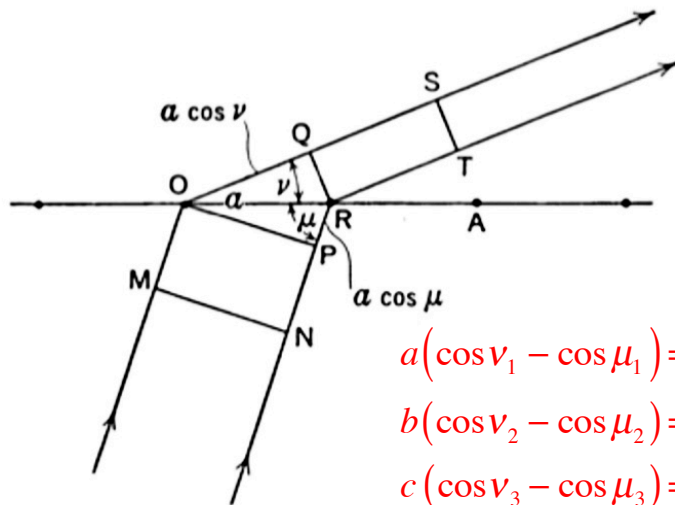
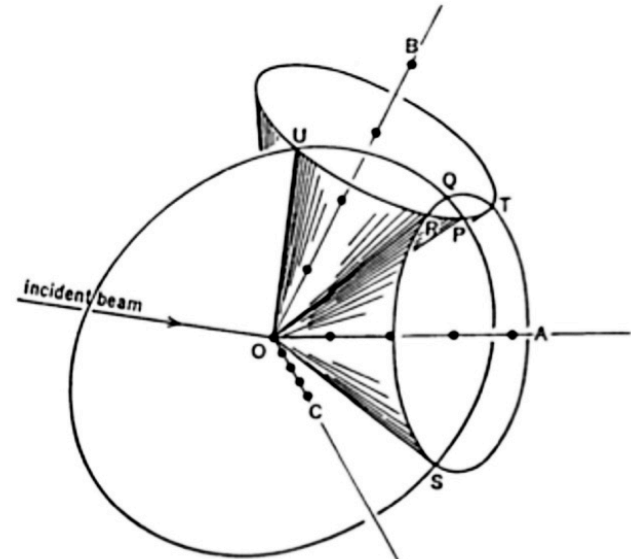
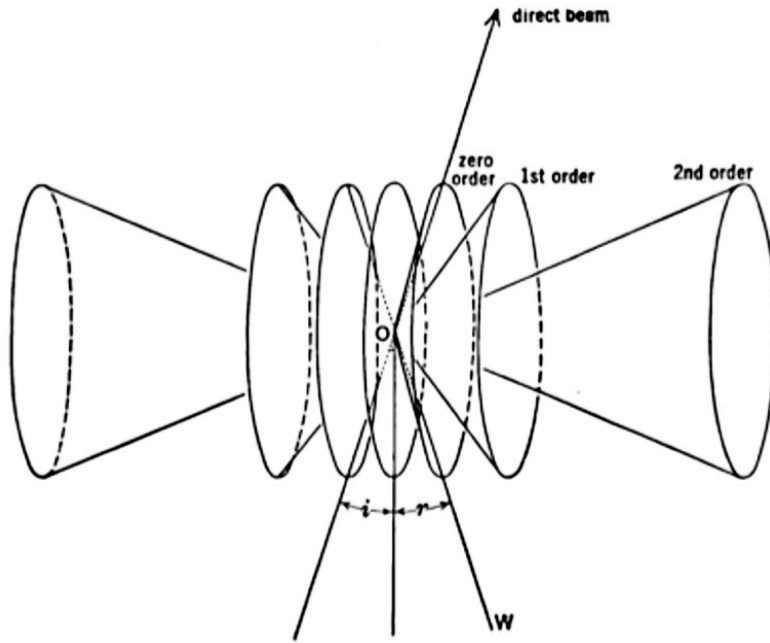
$$\vec{a} \cdot \hat{s} = |\vec{a}| \cos \nu$$

**Laue
diffraction
condition**

$$\vec{a} \cdot (\hat{s} - \hat{s}_0) = |\vec{a}| (\cos \nu - \cos \mu) = h\lambda$$



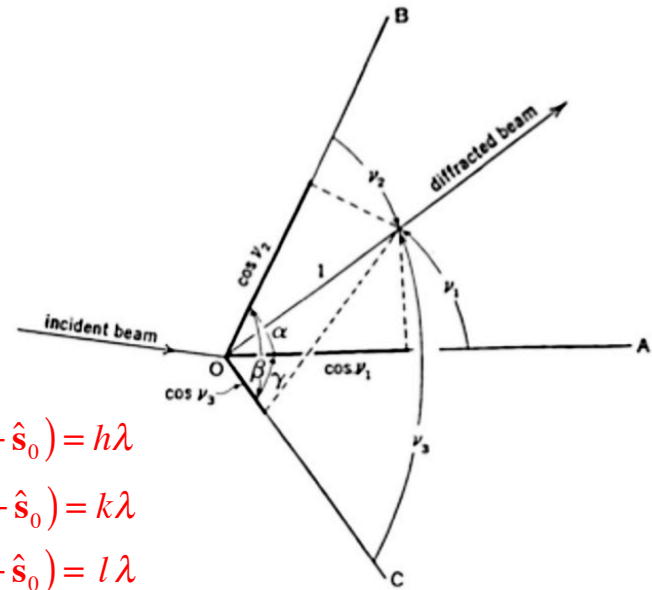
Laue diffraction by a three-dimensional lattice grating



$$a(\cos \nu_1 - \cos \mu_1) = \mathbf{a} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = h\lambda$$

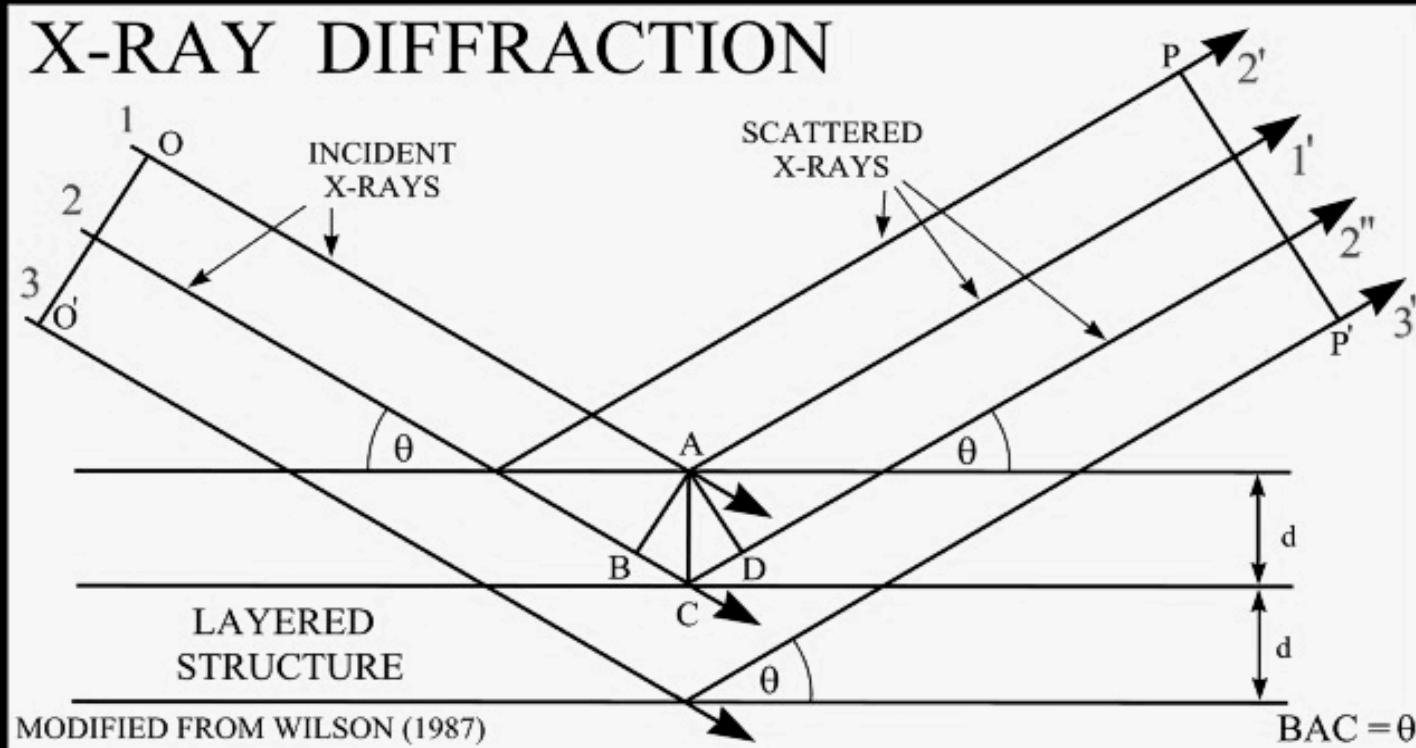
$$b(\cos \nu_2 - \cos \mu_2) = \mathbf{b} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = k\lambda$$

$$c(\cos \nu_3 - \cos \mu_3) = \mathbf{c} \cdot (\hat{\mathbf{s}} - \hat{\mathbf{s}}_0) = l\lambda$$



The Bragg Law

$$2d \sin \theta = n\lambda$$



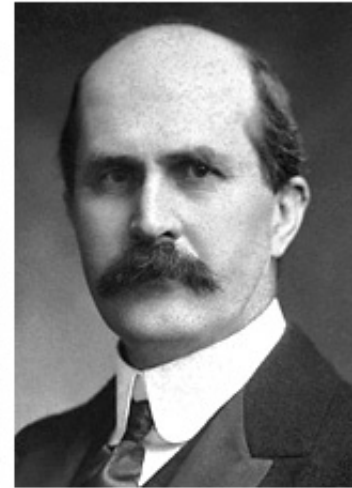
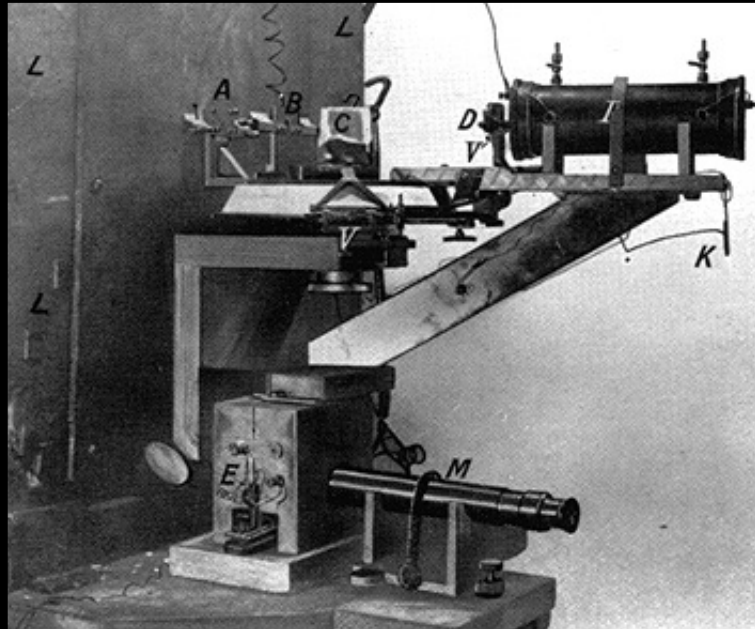
$$\angle ABC = \angle ADC = 90^\circ$$

$$\angle ACB = \angle ADB = 90^\circ - \theta$$

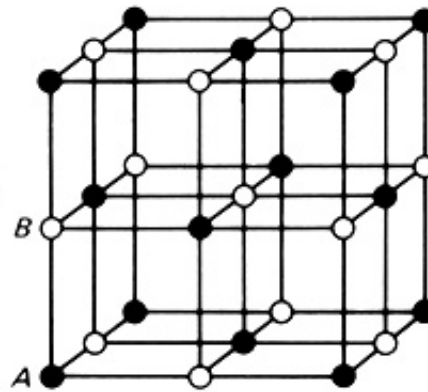
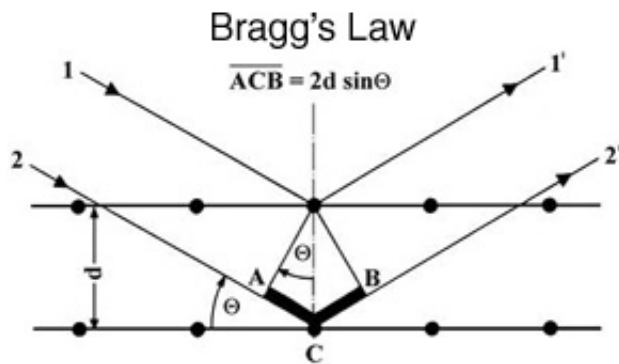
$$\angle BAC = \angle DAC = \theta$$

$$AC = d$$

$$BC = DC = d \sin \theta$$



William Henry Bragg

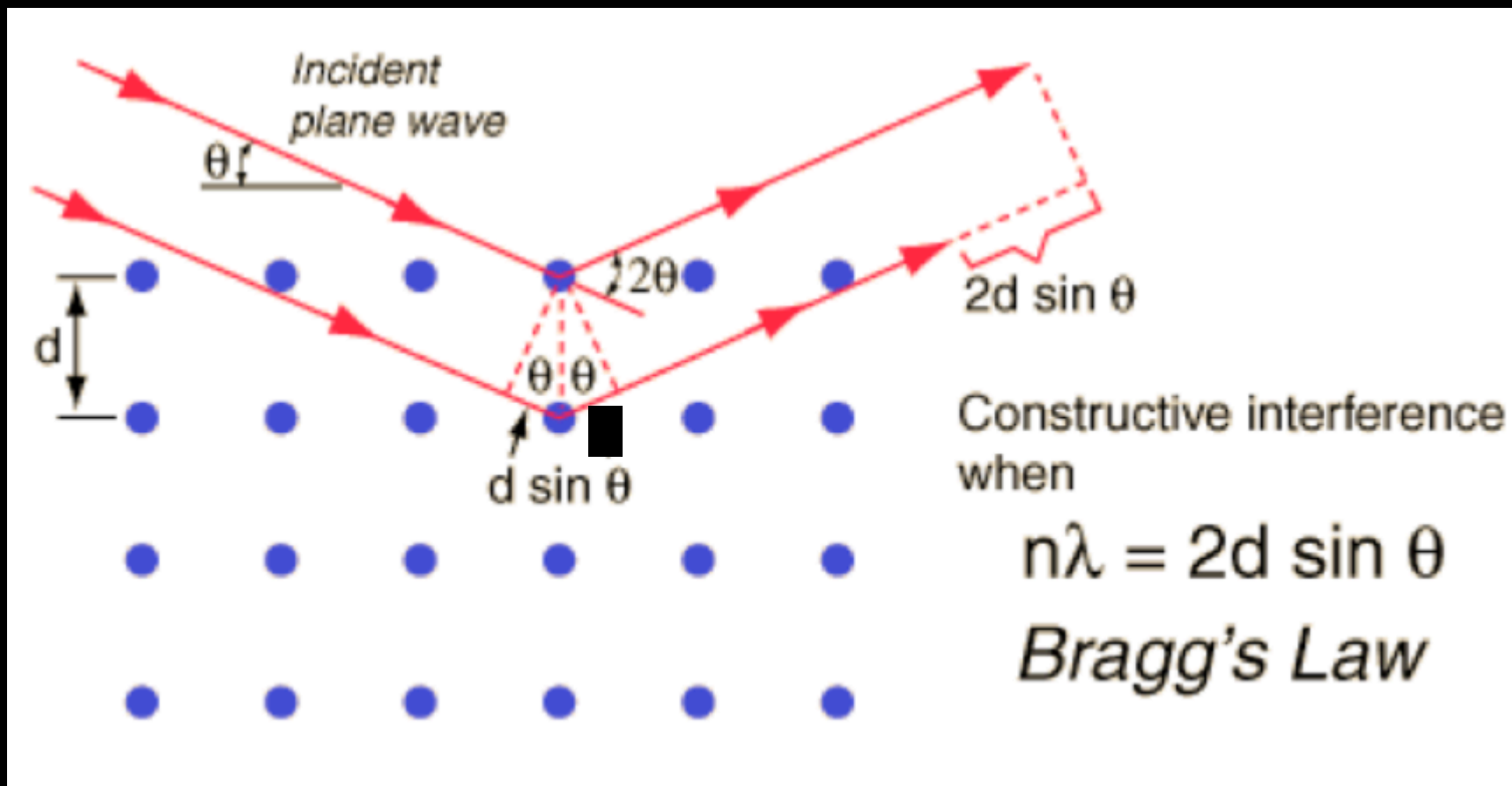


For NaCl
 $AB = 2.8 \times 10^{-8} \text{ cm}$



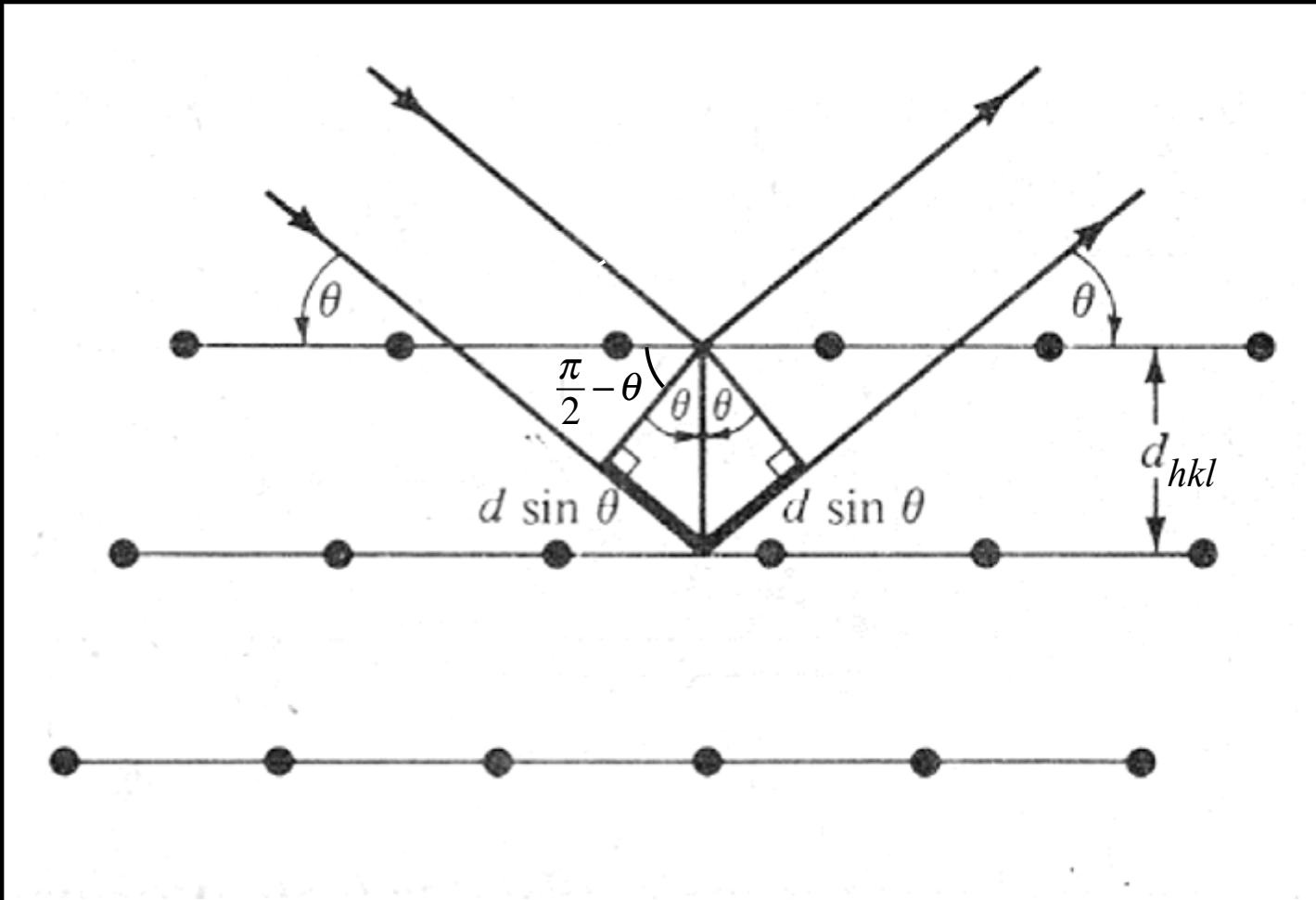
William Lawrence Bragg

Bragg's law describes constructive interference of X-ray waves scattered by a crystal as specular reflection from families of parallel crystal lattice planes



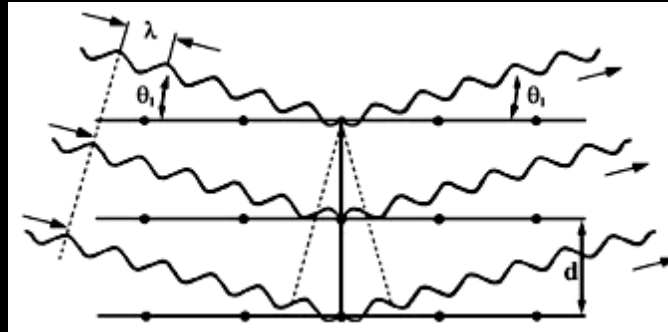
Bragg reflection from families of parallel hkl lattice planes

$$2d_{hkl} \sin \theta = n\lambda, \quad 2\left(\frac{d_{hkl}}{n}\right) \sin \theta = \lambda, \quad 2d_{nh\ nk\ n\ell} \sin \theta = \lambda$$



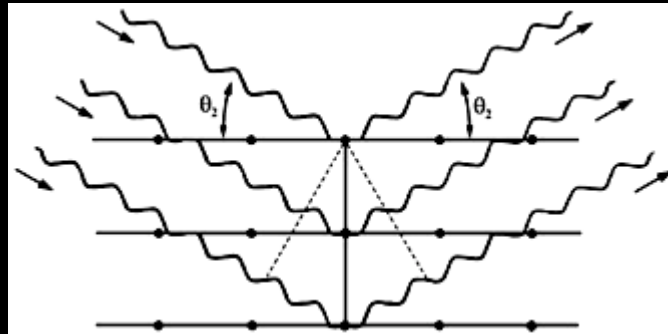
Orders of Bragg reflection from a family of parallel hkl lattice planes

1st order



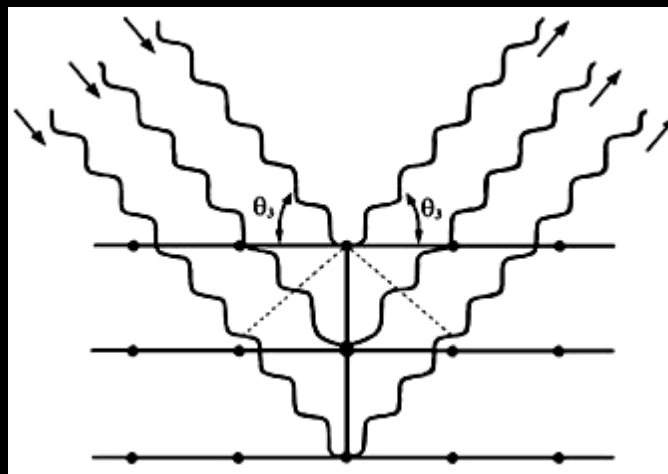
$$2d \sin \theta_1 = \lambda$$

2nd order



$$2d \sin \theta_2 = 2\lambda$$

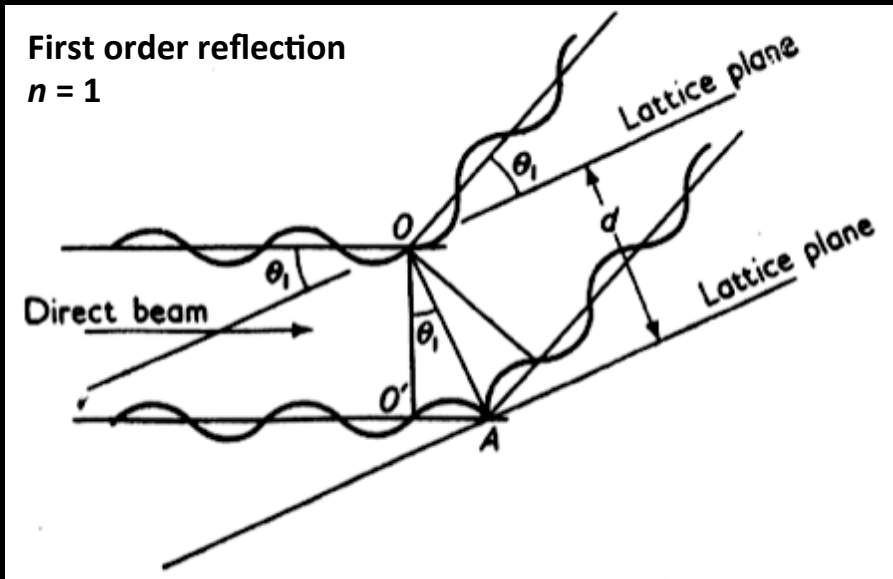
3rd order



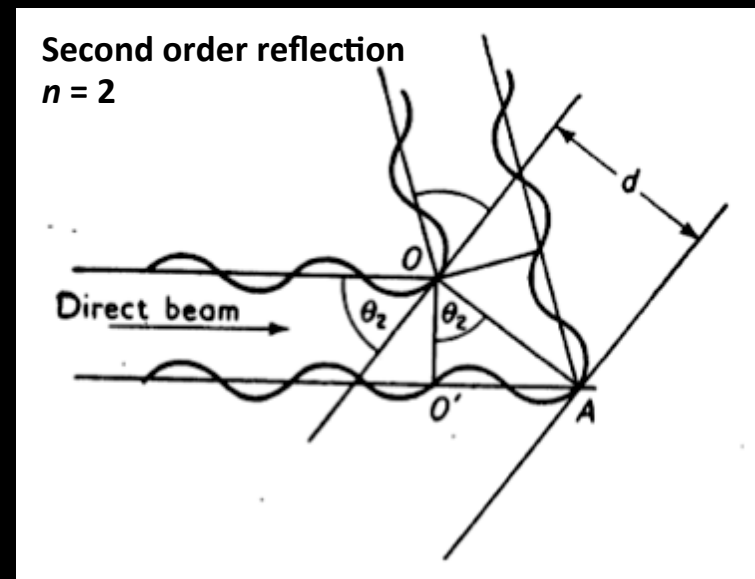
$$2d \sin \theta_3 = 3\lambda$$

Bragg reflections from a family of parallel hkl lattice planes

$$2d_{hkl} \sin \theta = n\lambda, \quad 2\left(\frac{d_{hkl}}{n}\right) \sin \theta = \lambda, \quad 2d_{nhnknl} \sin \theta = \lambda$$



$$2d_{hkl} \sin \theta_1 = \lambda$$



$$2d_{hkl} \sin \theta_2 = 2\lambda$$

$$2\left(\frac{d_{hkl}}{2}\right) \sin \theta_2 = \lambda$$

$$2d_{2h\ 2k\ 2l} \sin \theta_2 = \lambda$$

Monochromatic X-ray diffraction

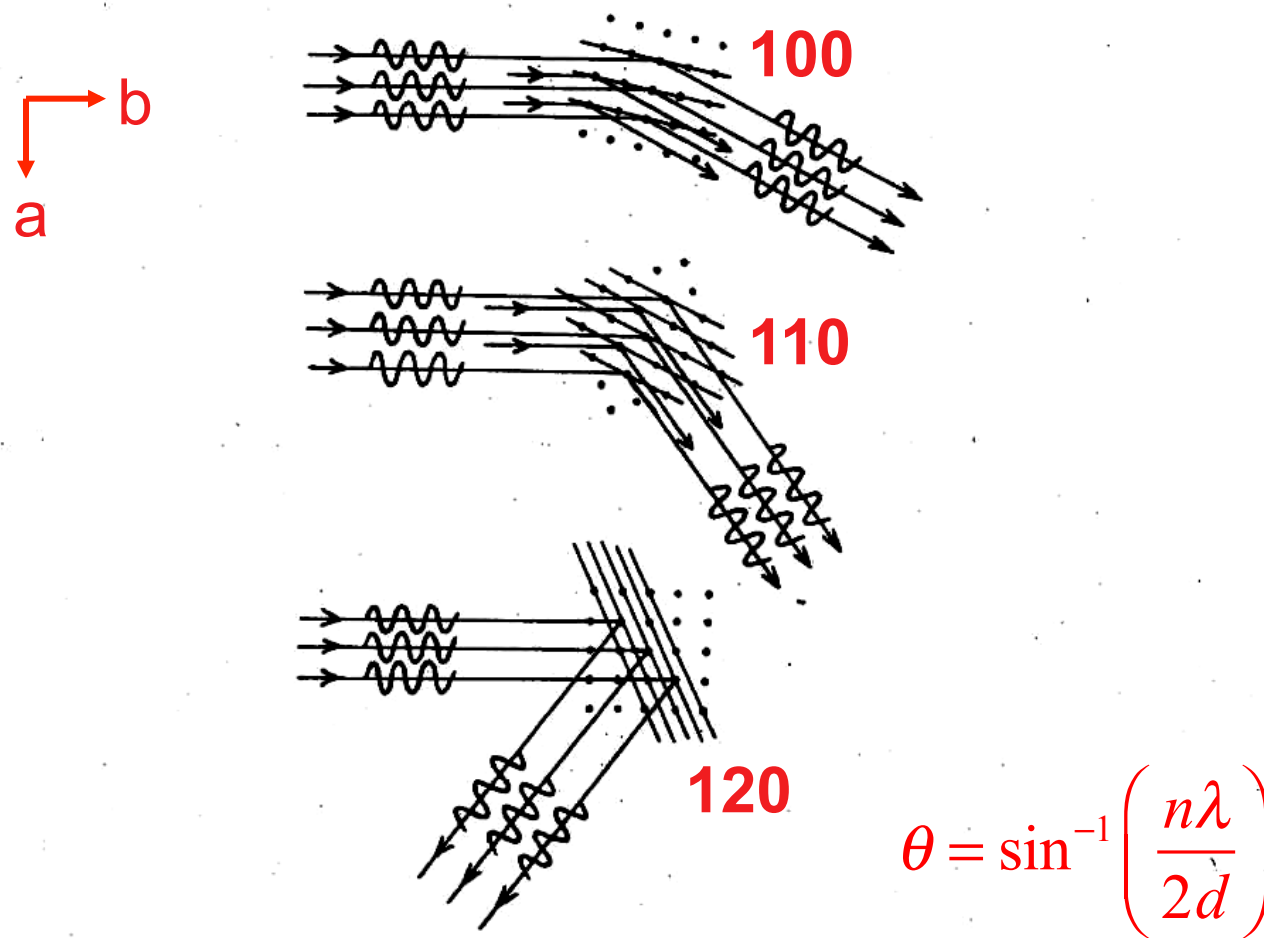


FIG. 113. When a crystal is rotated in a monochromatic X-ray beam, the various planes of atoms are able to reflect only at particular angles; the angle at which reflection takes place depends on the spacing of the planes. The more widely spaced planes give reflections near the primary beam; the more closely spaced planes reflect at the larger angles.

Polychromatic X-ray diffraction

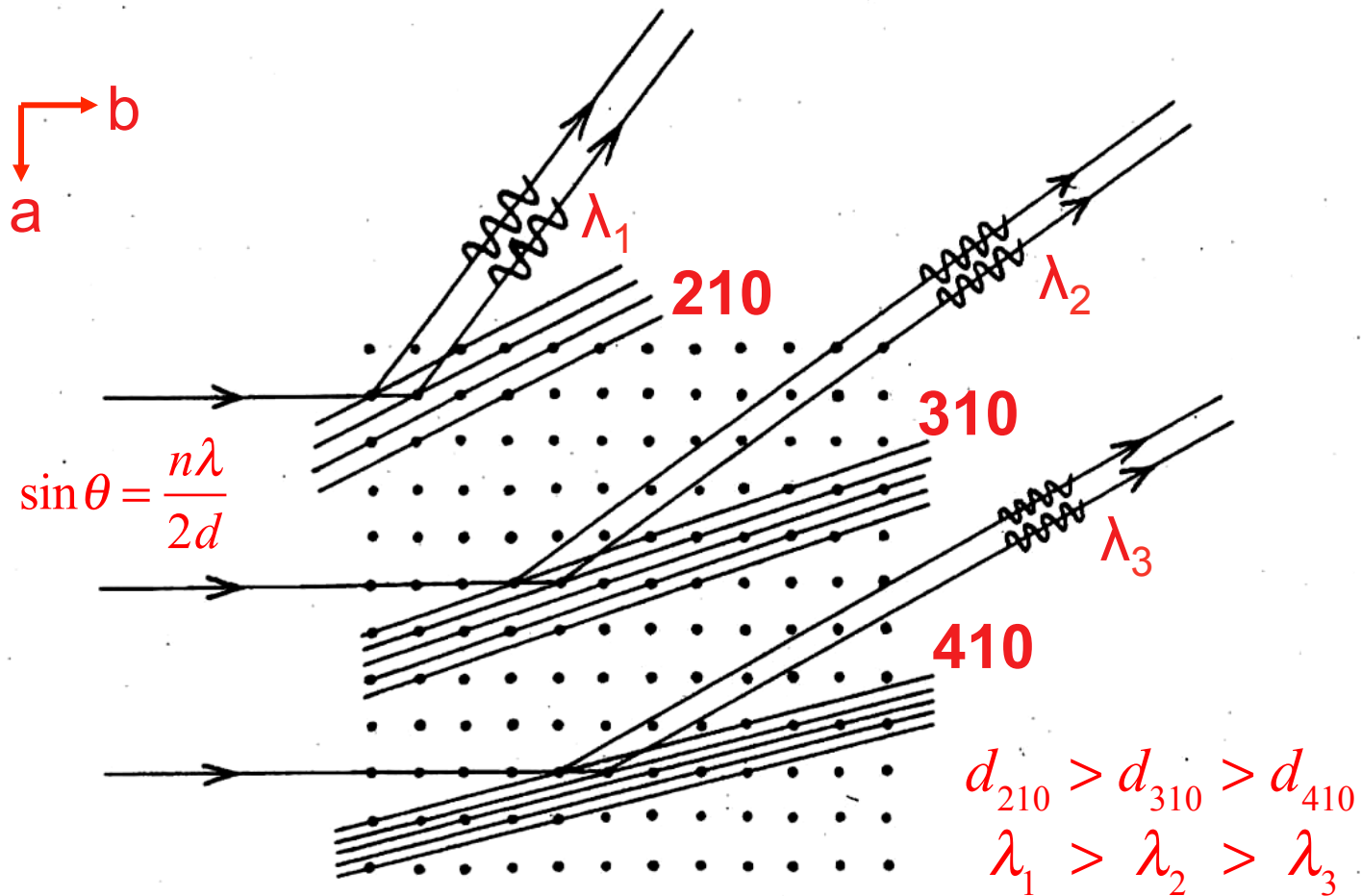


FIG. 111. In Laue's experiment, each spot on the plate is due to the reflection of X-rays by a different set of atomic planes acting as X-ray mirrors. The more widely spaced planes reflect the longer waves in the beam, while the more closely spaced planes reflect the shorter waves, at smaller angles.

Polychromatic X-ray diffraction

According to the Bragg law $2d_{hkl} \sin \theta_{hkl} = \lambda$, the discrete lattice spacings d_{hkl} select discrete monochromatic wavelengths $\lambda_i = 2d_{hkl} (\sin \theta_{hkl})$ from a polychromatic X-ray beam.

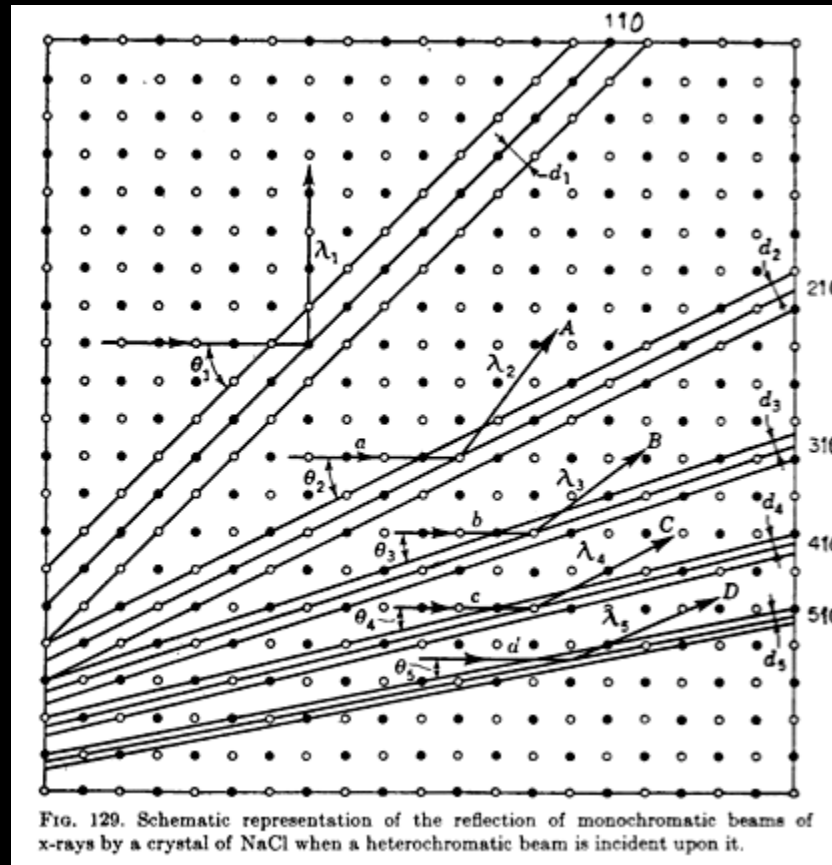


FIG. 129. Schematic representation of the reflection of monochromatic beams of x-rays by a crystal of NaCl when a heterochromatic beam is incident upon it.

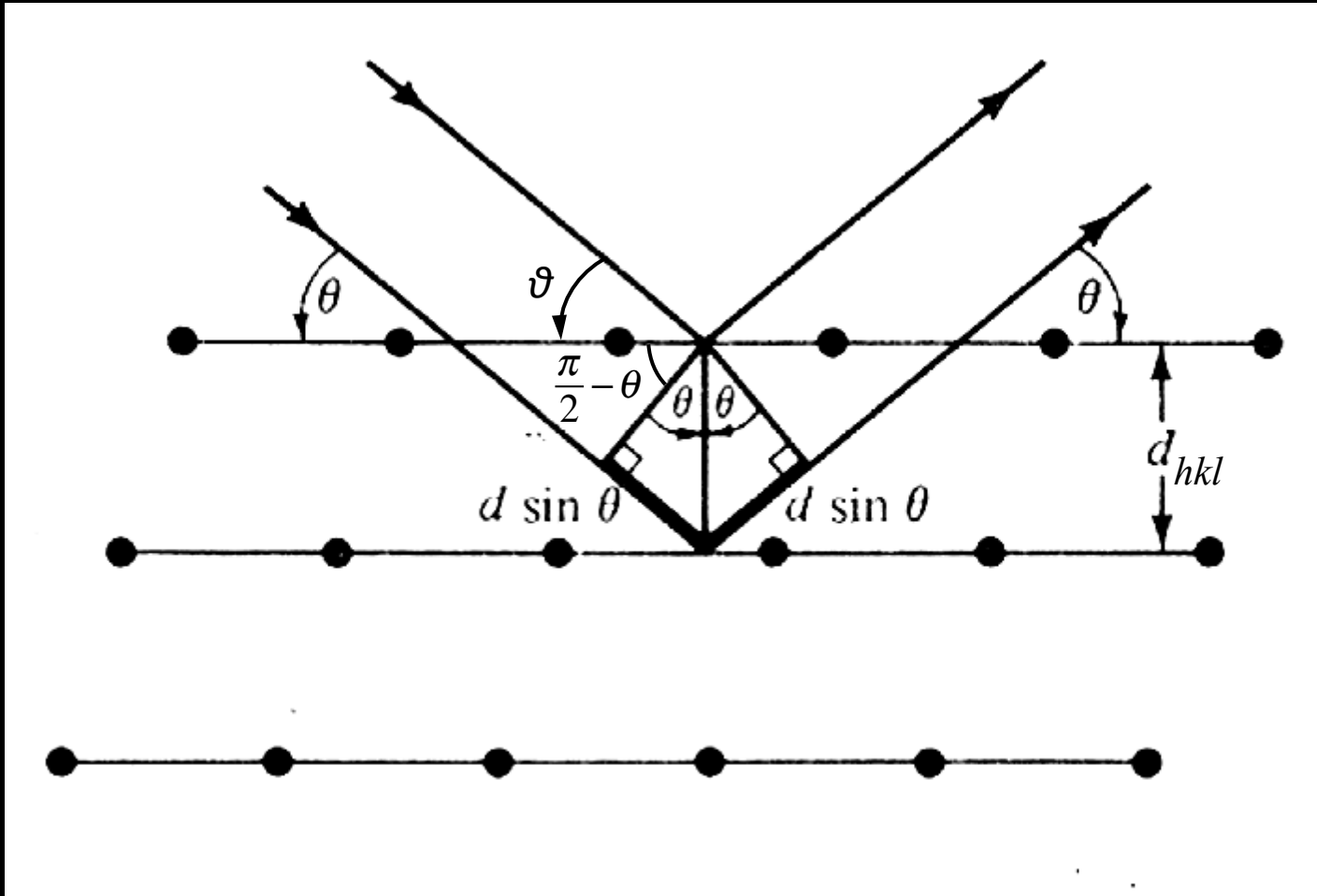
$$d_{110} > d_{210} > d_{310} > d_{410} > d_{510}$$

$$\theta_1 > \theta_2 > \theta_3 > \theta_4 > \theta_5$$

$$\lambda_1 > \lambda_2 > \lambda_3 > \lambda_4 > \lambda_5$$

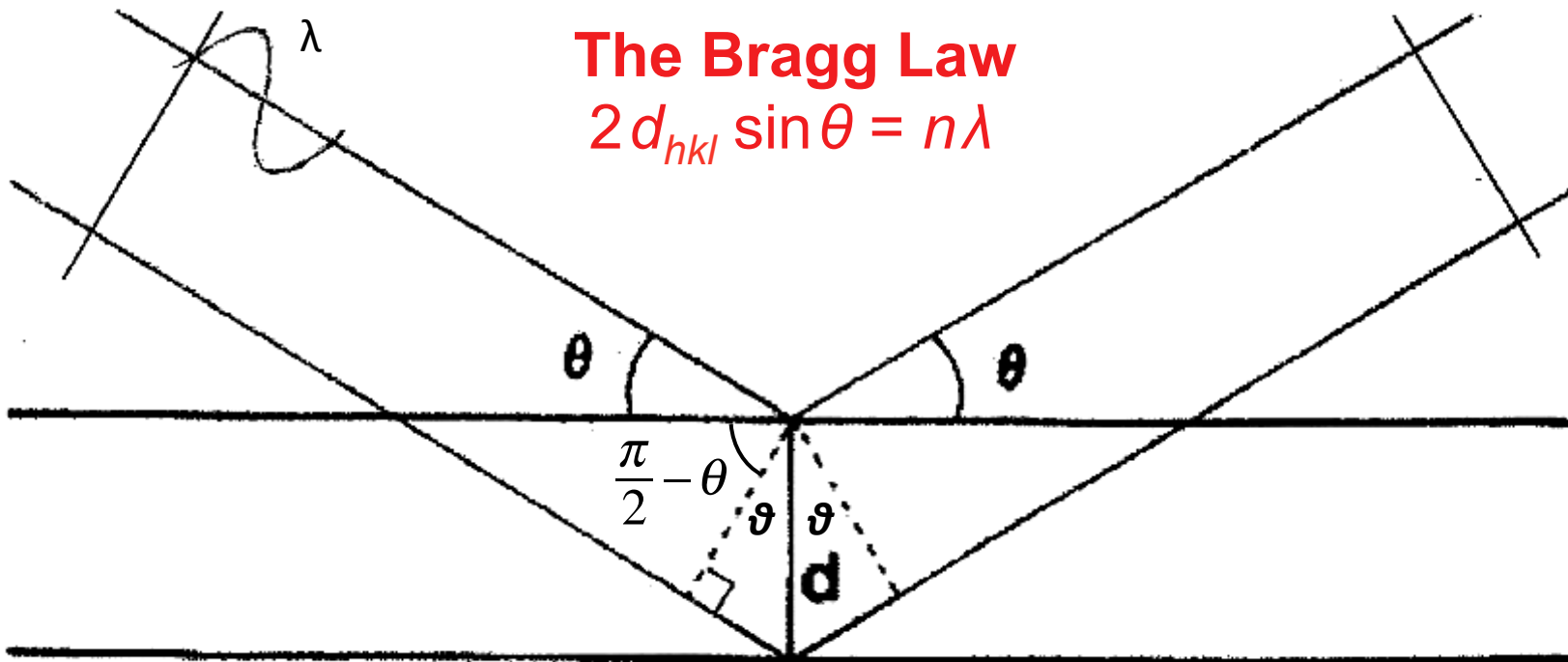
Bragg reflection from families of parallel hkl lattice planes

$$2d_{hkl} \sin \theta = n\lambda, \quad 2\left(\frac{d_{hkl}}{n}\right) \sin \theta = \lambda, \quad 2d_{nhnknl} \sin \theta = \lambda$$

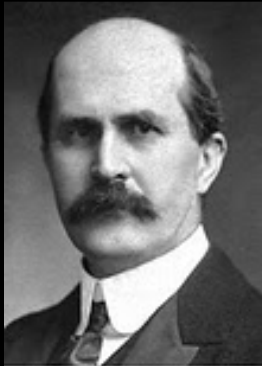


The Bragg Law

$$2d_{hkl} \sin \theta = n\lambda$$



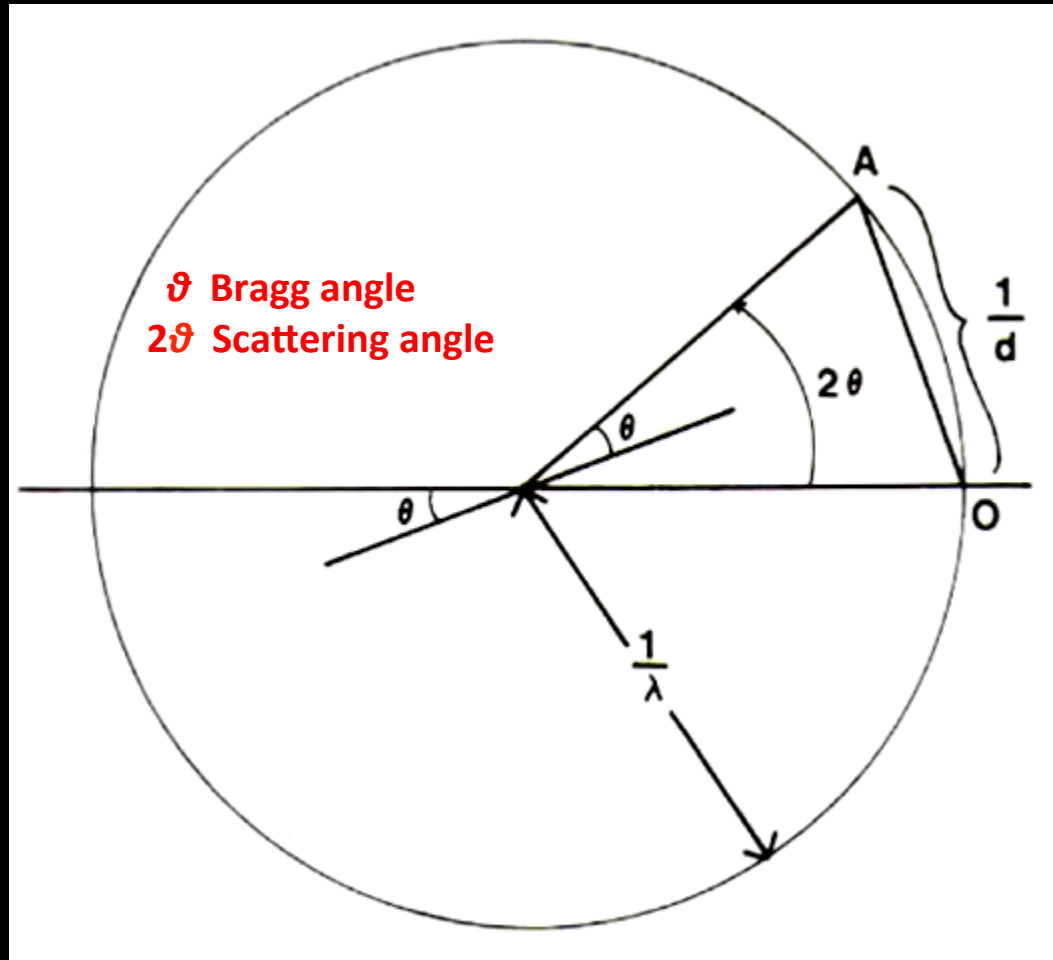
The Ewald construction for the Bragg reflection law



Sir William Henry Bragg
(1862 - 1942)



William Lawrence Bragg
(1890 - 1971)



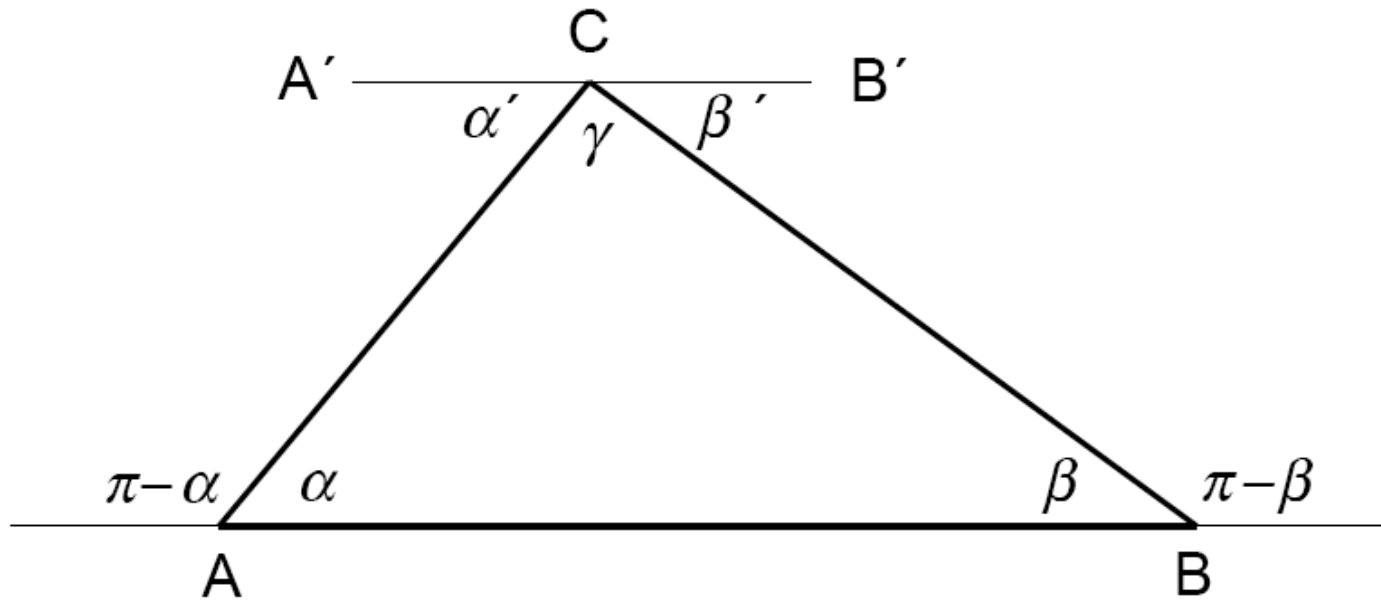
Paul Peter Ewald
1888-1985

{	Wm. Henry & Wm. Lawrence Bragg (1913)	$2d \sin \theta = \lambda$
	Paul Peter Ewald (1913)	$\sin \theta = \frac{1}{2} \frac{\lambda}{d} = \left(\frac{1}{d} \right) / \left(\frac{2}{\lambda} \right)$

Robert M. Sweet (1985). Introduction to Crystallography. In *Methods in Enzymology*, v.114.

<http://sandwalk.blogspot.com/2008/03/nobel-laureates-sir-william-henry-bragg.html>

<http://wwwuser.gwdg.de/~nprofwg/ForeignScientists.html>



$$A'B' \parallel AB$$

$$\alpha' = \alpha$$

$$\beta' = \beta$$

$$\alpha' + \gamma + \beta' = \alpha + \beta + \gamma = \pi$$

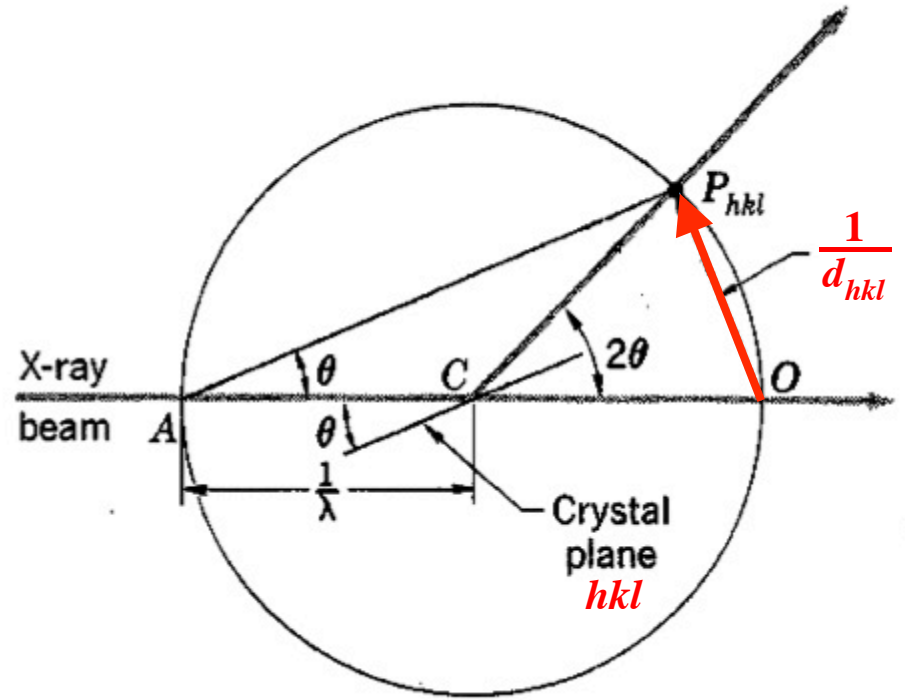
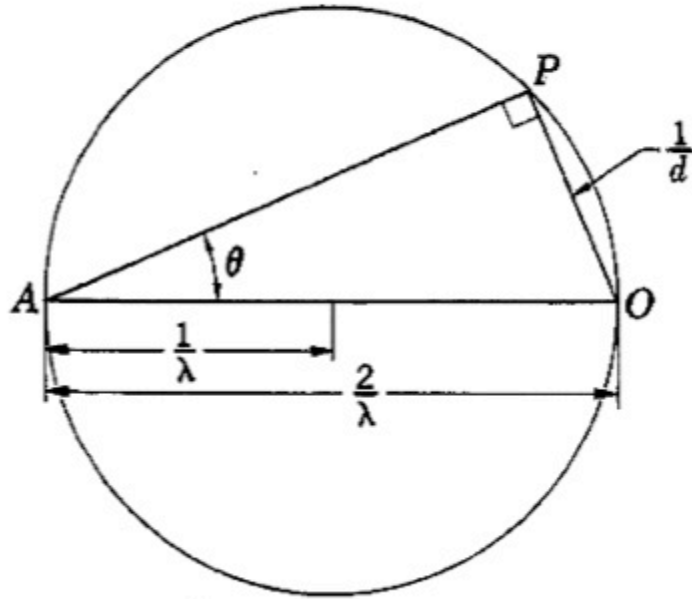
$$\alpha + \beta + (\alpha + \beta) = \pi$$

$$2\alpha + 2\beta = \pi$$

$$\alpha + \beta = \frac{\pi}{2}$$

The Ewald construction for the Bragg reflection law

ϑ Bragg angle
 2ϑ Scattering angle



{	Wm. Henry & Wm. Lawrence Bragg (1913)	$2d \sin \theta = \lambda$
	Paul Peter Ewald (1913)	$\sin \theta = \frac{1}{d} \frac{\lambda}{2} = \left(\frac{1}{d} \right) / \left(\frac{2}{\lambda} \right)$

The Ewald construction for the Bragg reflection law

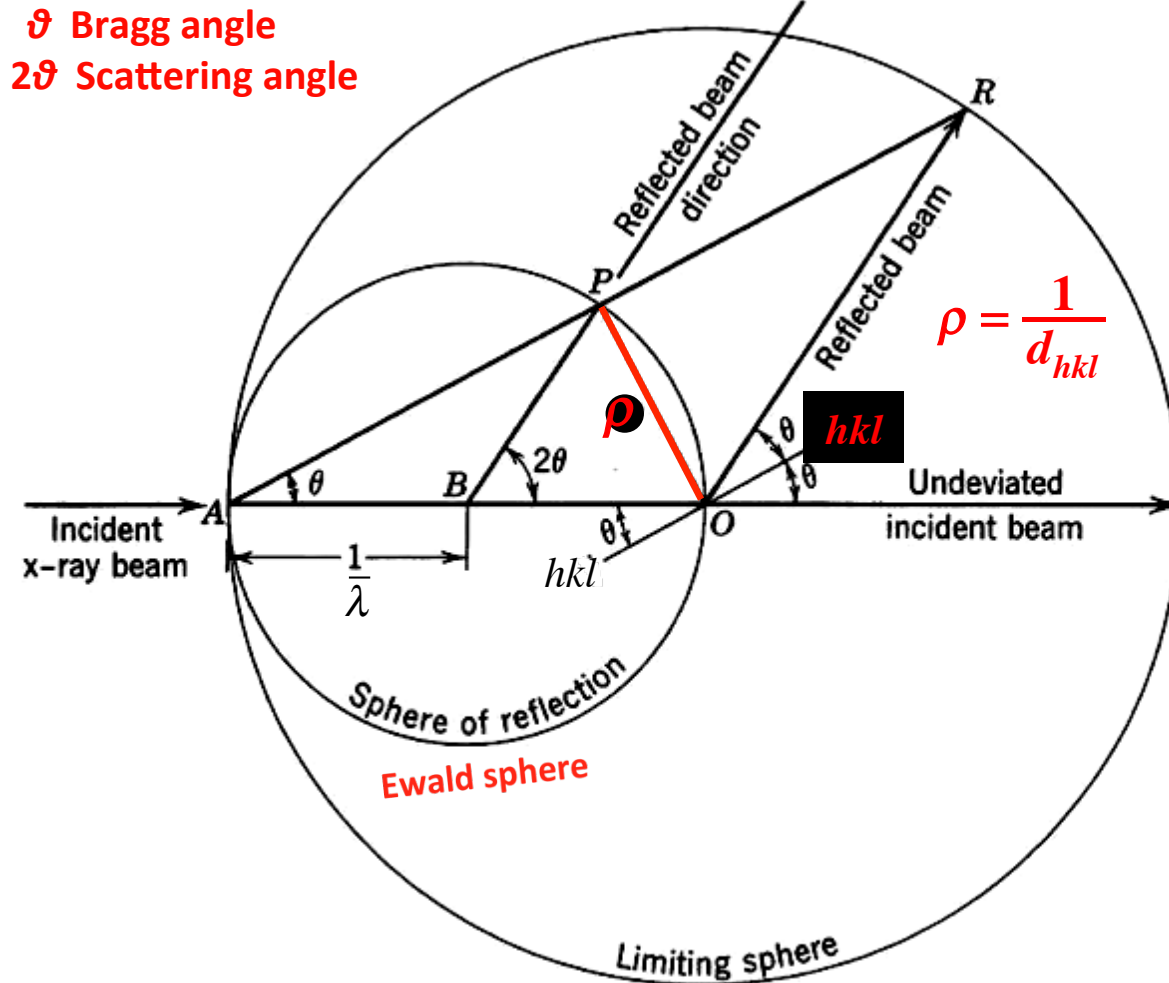
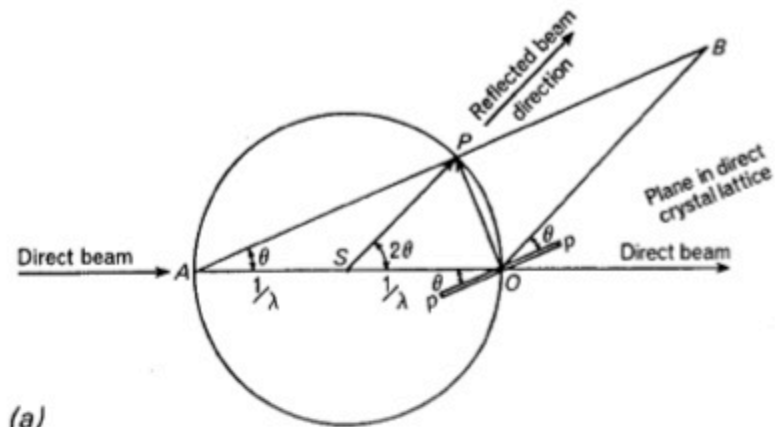
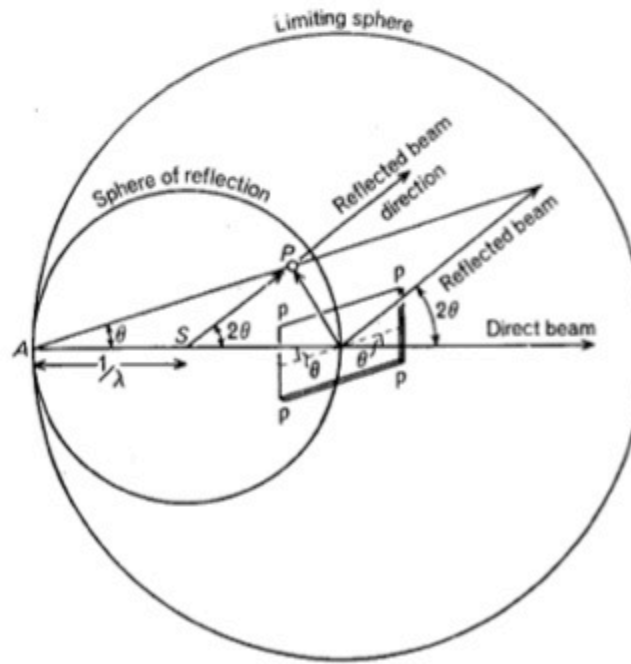


Fig. 3-8. The geometrical conditions for reflection in reciprocal space.
(Courtesy of Bernal, *Proc. Roy. Soc. (London)*, 113A, 117.)



(a)



(b)

FIG. 3. The Ewald construction showing the condition for reflection in (a) two dimensions, and (b) three dimensions.

The Ewald construction

$$2d_{hkl} \sin \theta_{hkl} = \lambda$$

$$\sin \theta_{hkl} = \frac{\lambda}{2d_{hkl}} = \frac{1}{2} \frac{1/d_{hkl}}{1/\lambda} = \frac{d_{hkl}^*}{2/\lambda}$$

$$\mathbf{H} = \frac{\hat{\mathbf{s}} - \hat{\mathbf{s}}_0}{\lambda}$$

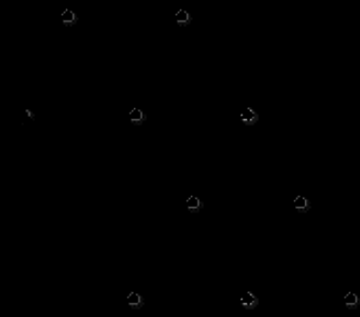
$$|\mathbf{H}| = d_{hkl}^* = 2 \left(\frac{\sin \theta_{hkl}}{\lambda} \right)$$

$$\angle \text{AOX} = \pi - 2\theta$$

$$\therefore \angle \text{OAX} = \angle \text{OXA} = \theta$$

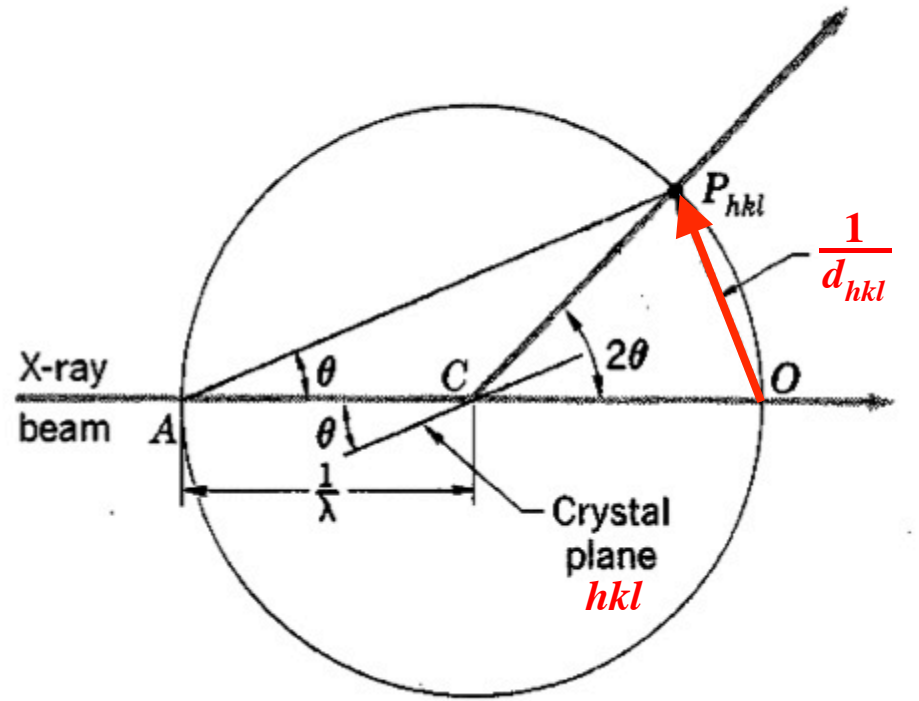
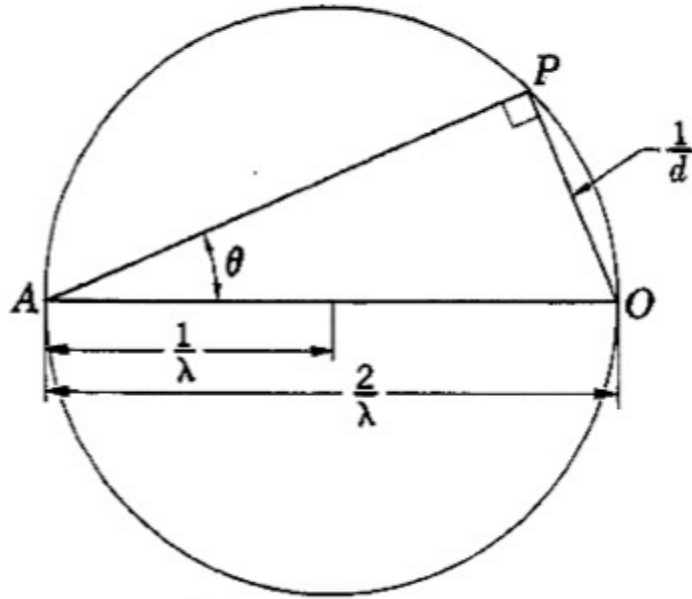
$$\angle \text{OXY} = \angle \text{OYX} = \frac{1}{2}(\pi - 2\theta) = \frac{\pi}{2} - \theta$$

$$\therefore \angle \text{AXY} = \frac{\pi}{2}$$



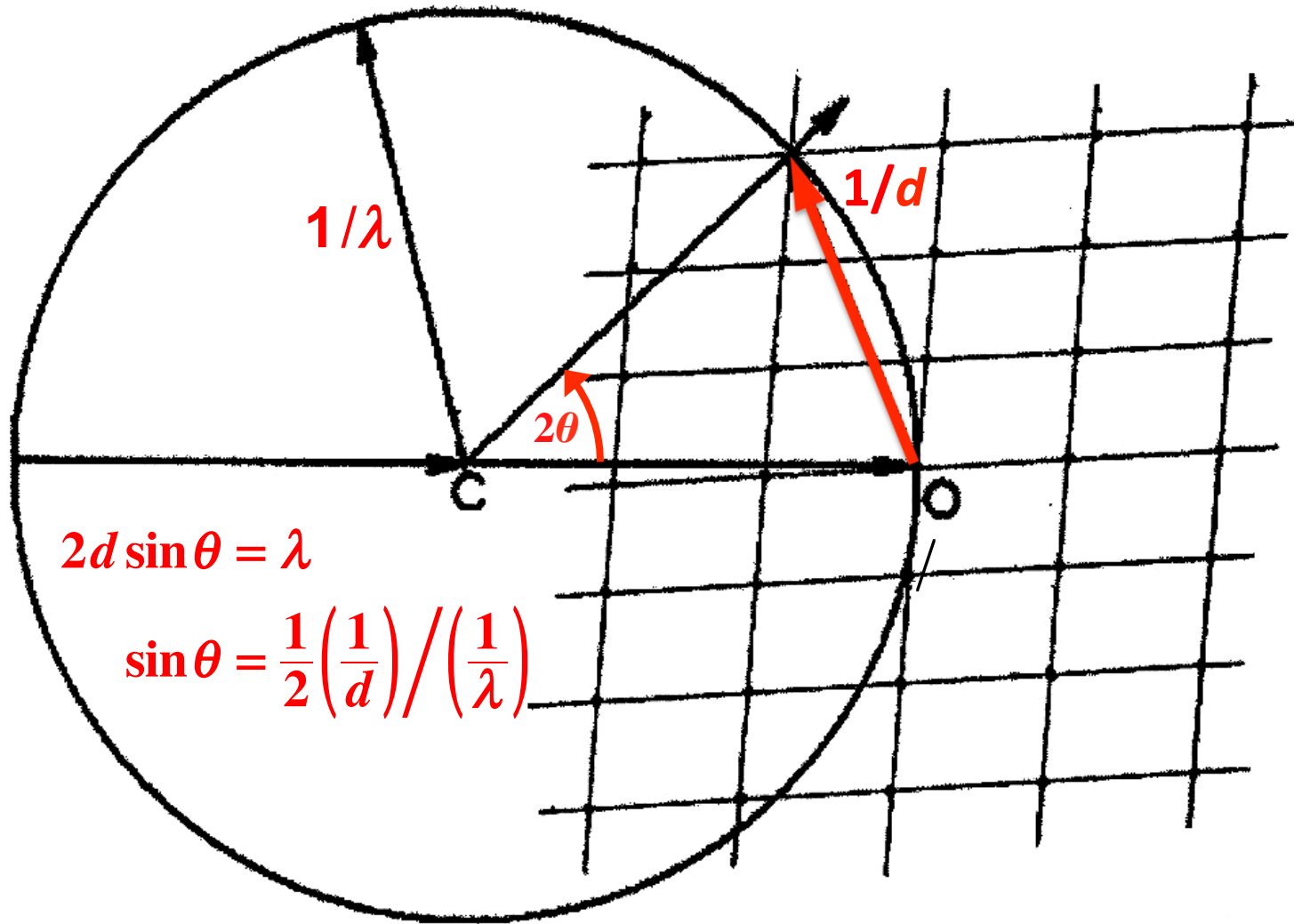
The Ewald construction for the Bragg reflection law

ϑ Bragg angle
 2ϑ Scattering angle



{	Wm. Henry & Wm. Lawrence Bragg (1913)	$2d \sin \theta = \lambda$
	Paul Peter Ewald (1913)	$\sin \theta = \frac{1}{d} \frac{\lambda}{2} = \left(\frac{1}{d} \right) / \left(\frac{2}{\lambda} \right)$

The Ewald construction on the reciprocal lattice



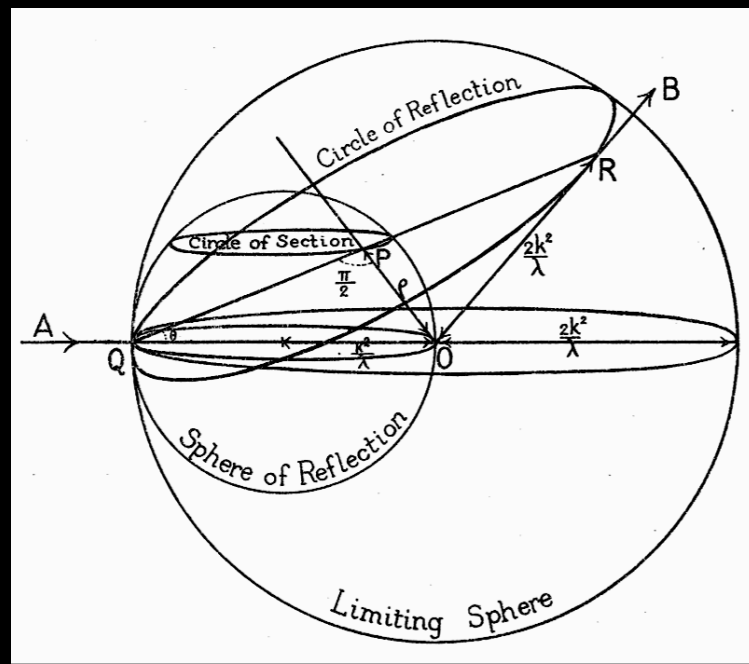
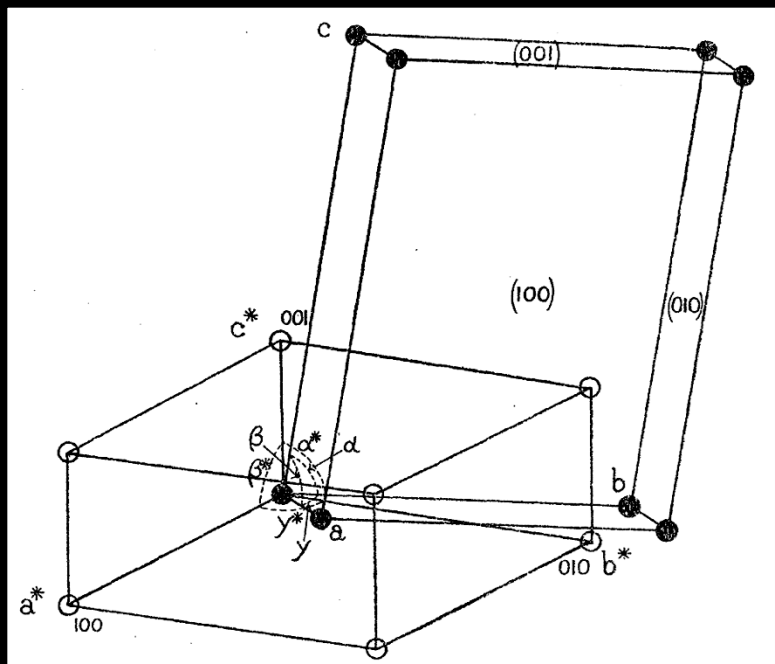


On the Interpretation of X-Ray, Single Crystal, Rotation Photographs.

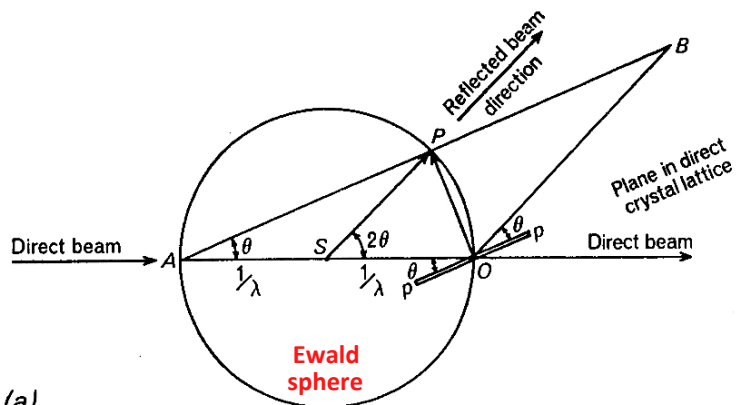
By J. D. BERNAL, B.A.

(Communicated by Sir William Bragg, F.R.S.—Received July 24, 1926.)

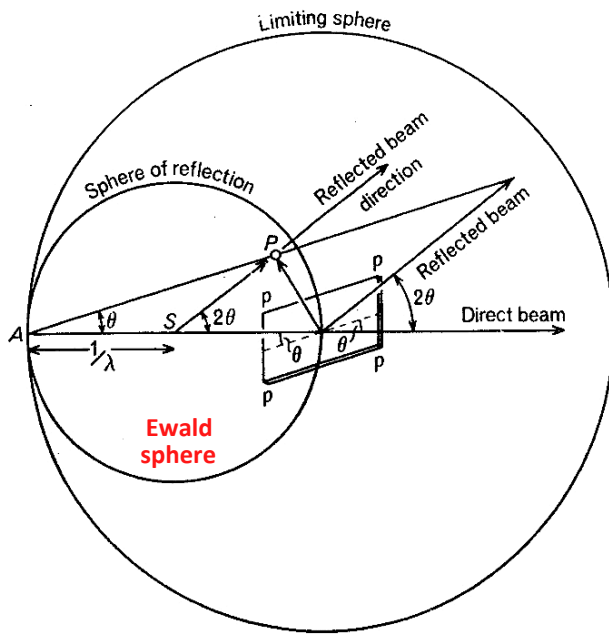
In the development of the study of crystals by X-rays the methods used divide themselves naturally into four types: the Bragg Ionisation Spectrometer method, the Laue method, the Powder method of Debye and Scherrer, and the Rotating Crystal method of Rinne, Schiebold and Polyani. The techniques of the first three of these methods are fully explained in such books as 'X-Rays and Crystal Structure,' by W. H. and W. L. Bragg, 'The Structure of Crystals,' by Wyckoff, and 'Krystalle und Rontgenstrahlen,' by Ewald, as well as in original papers. On the other hand, the rotation method is only



The Ewald sphere along side an Escher sphere

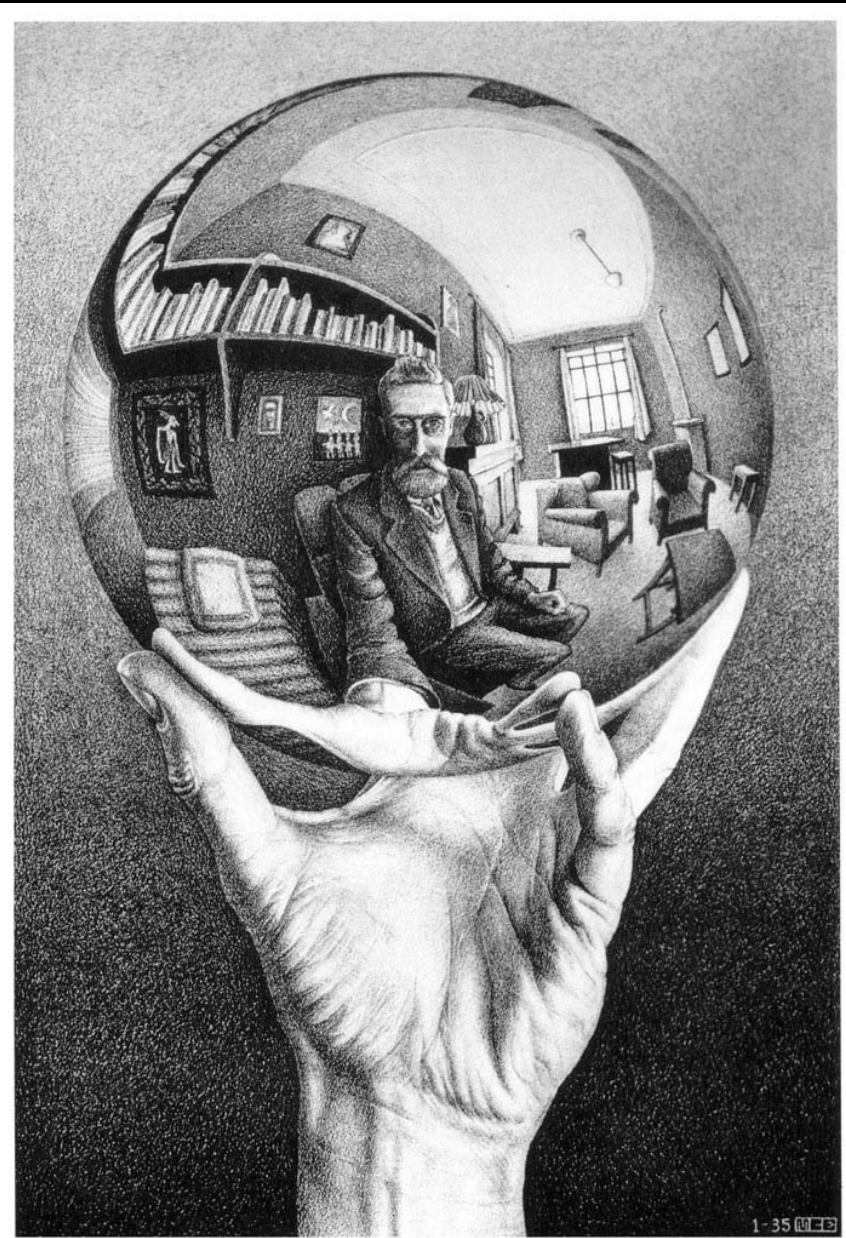


(a)

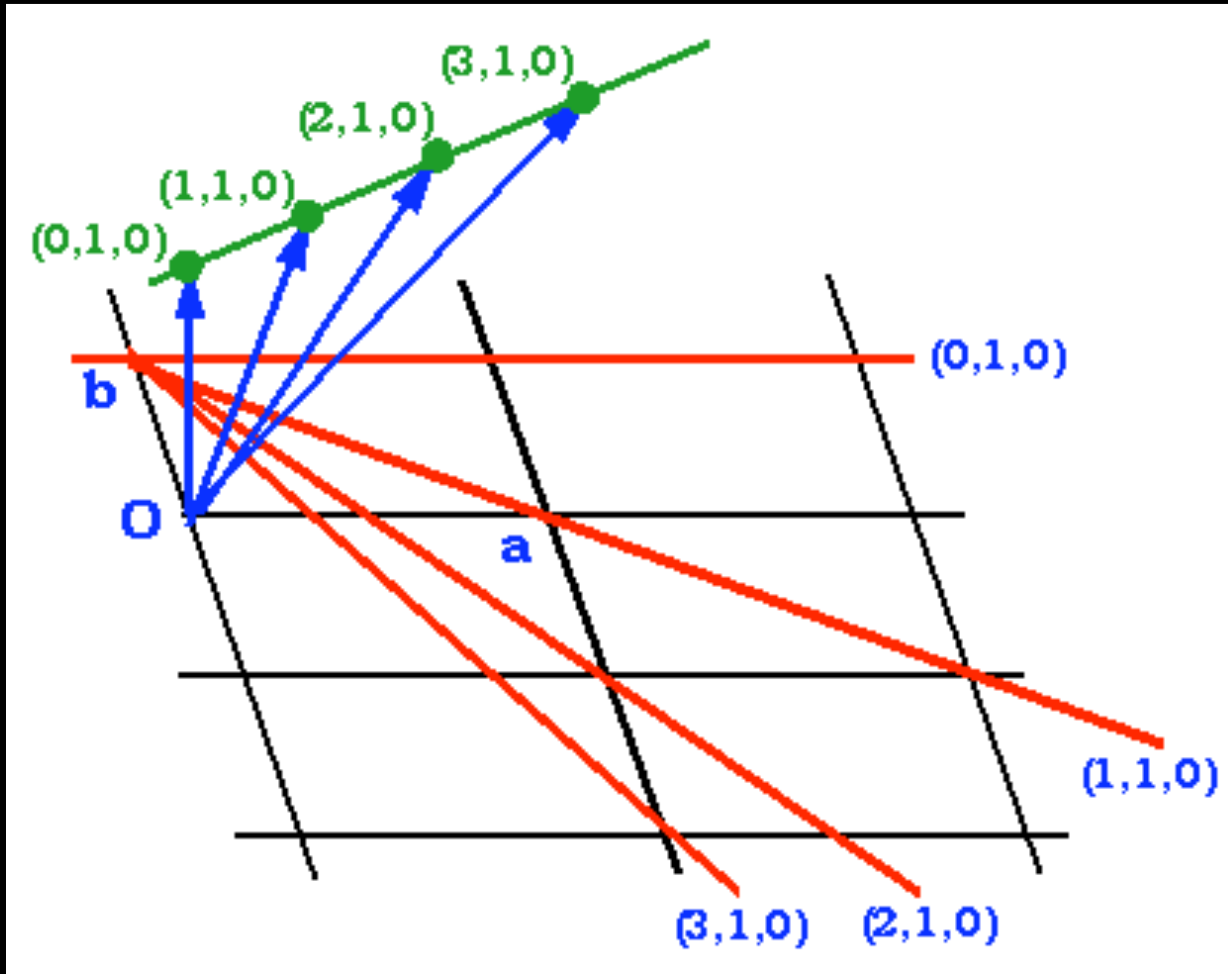


(b)

FIG. 3. The Ewald construction showing the condition for reflection in (a) two dimensions, and (b) three dimensions.

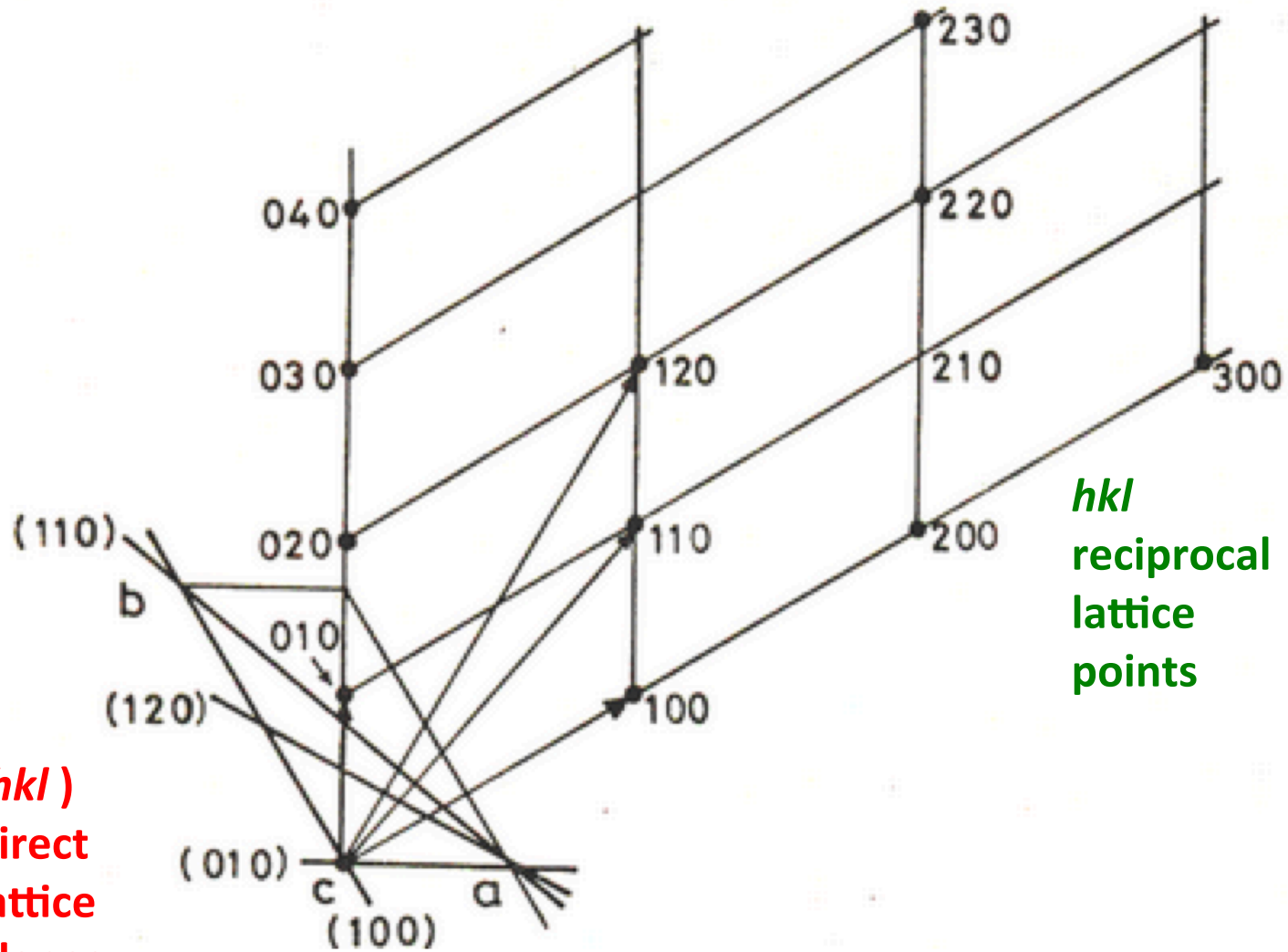


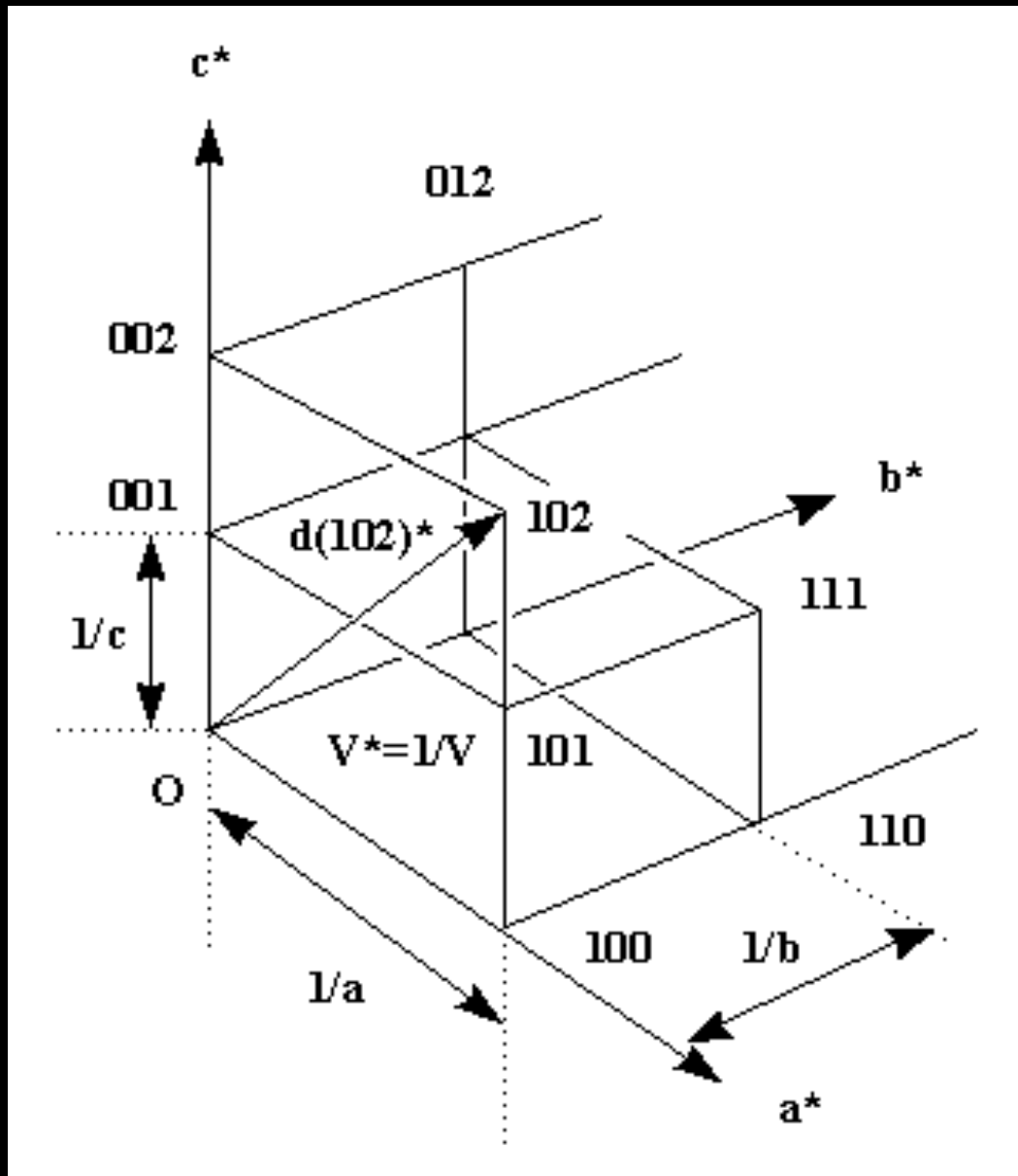
Reciprocal lattice points
Reciprocal lattice vectors

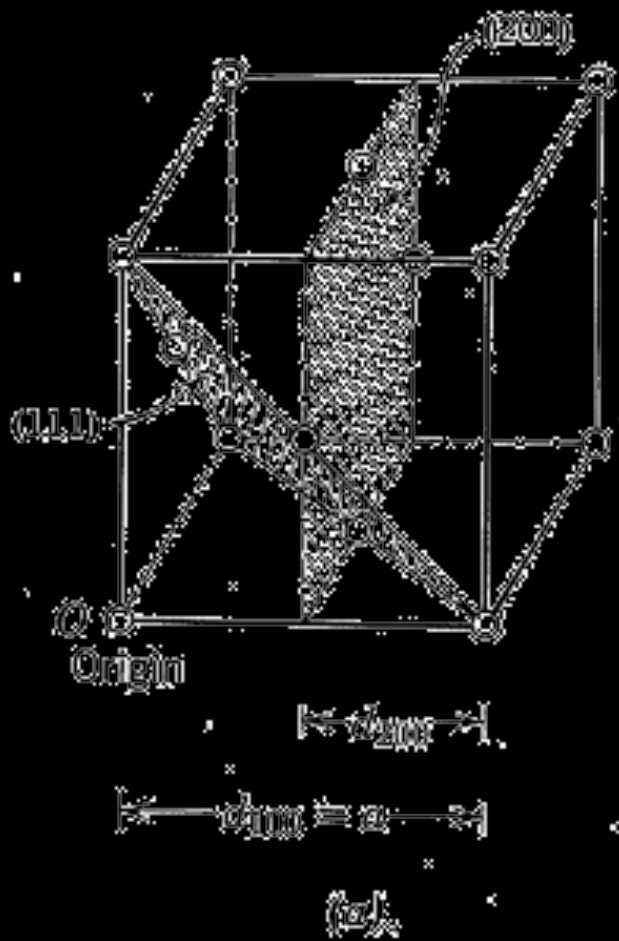


Direct lattice planes

Geometrical construction of a reciprocal lattice







a b c Direct Lattice and **a* b* c*** Reciprocal Lattice

N.B. In all crystal systems:

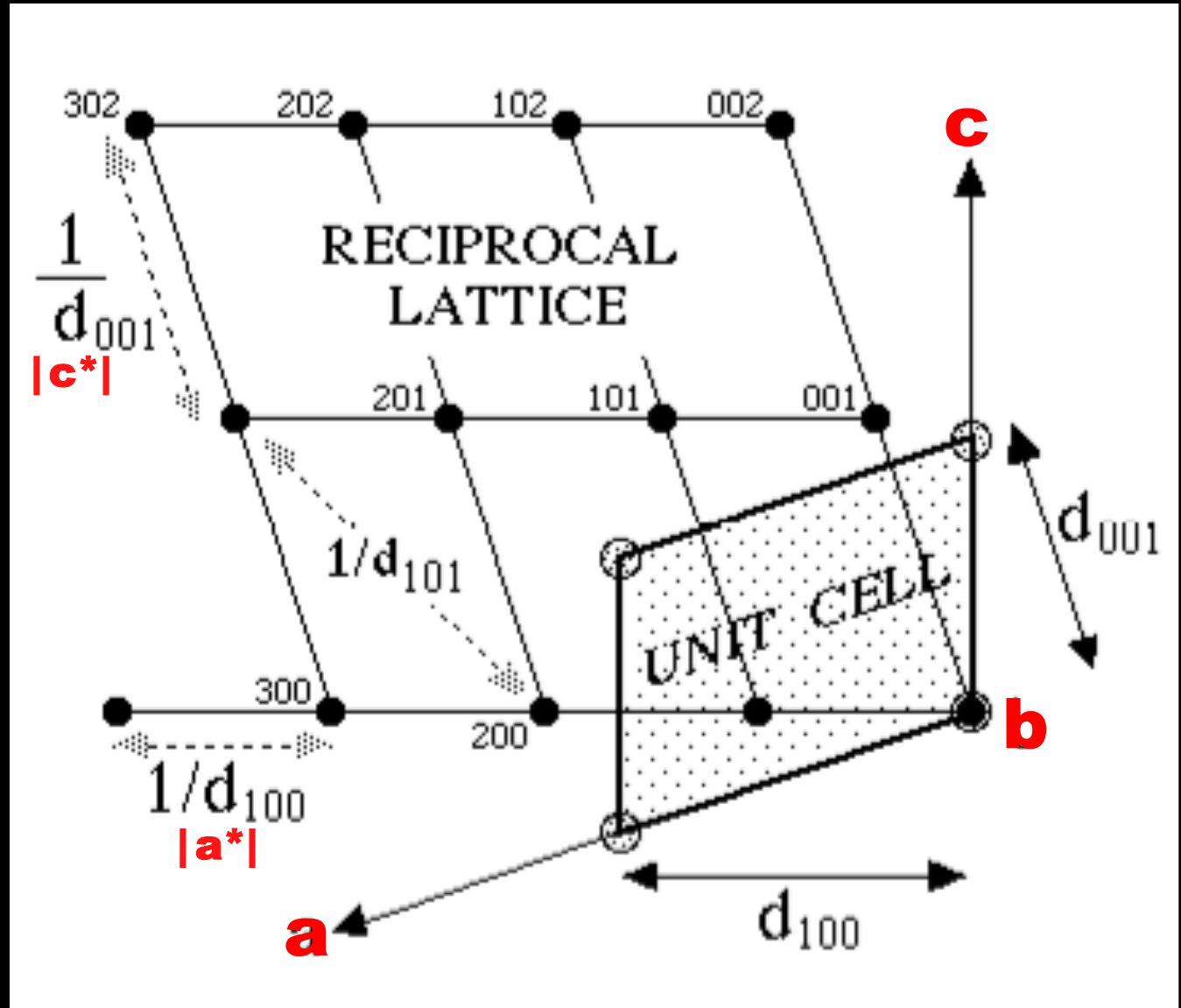
$$\left\{ \begin{array}{l} \mathbf{a}^* \perp \mathbf{b} \text{ and } \mathbf{c} \\ \mathbf{b}^* \perp \mathbf{c} \text{ and } \mathbf{a} \\ \mathbf{c}^* \perp \mathbf{a} \text{ and } \mathbf{b} \end{array} \right.$$

and

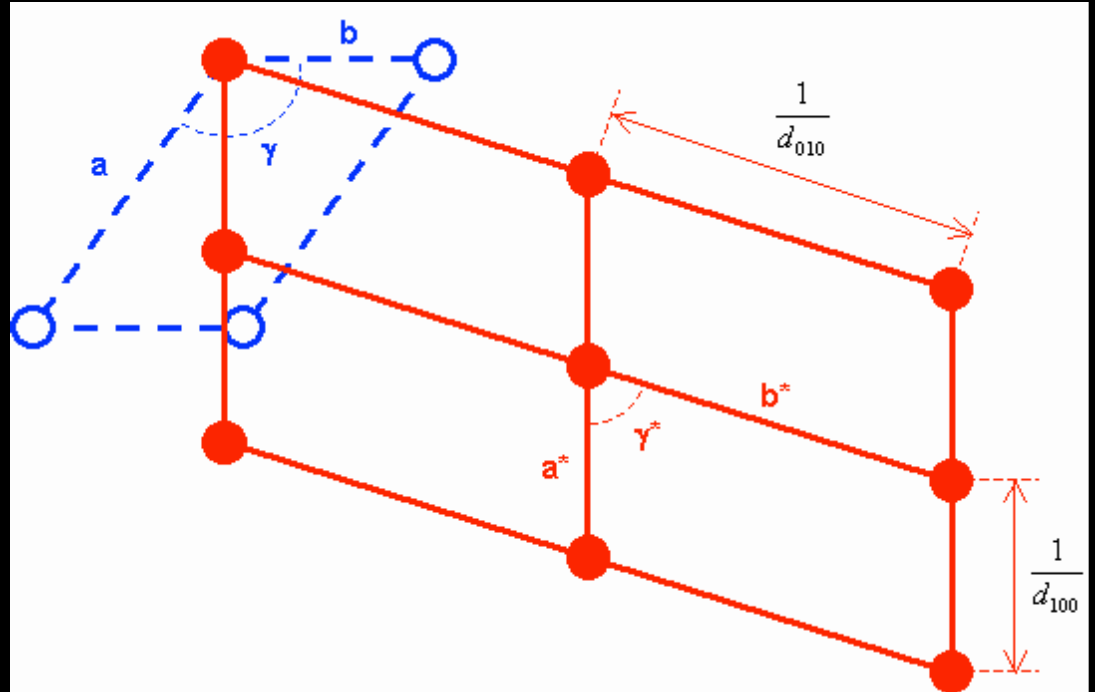
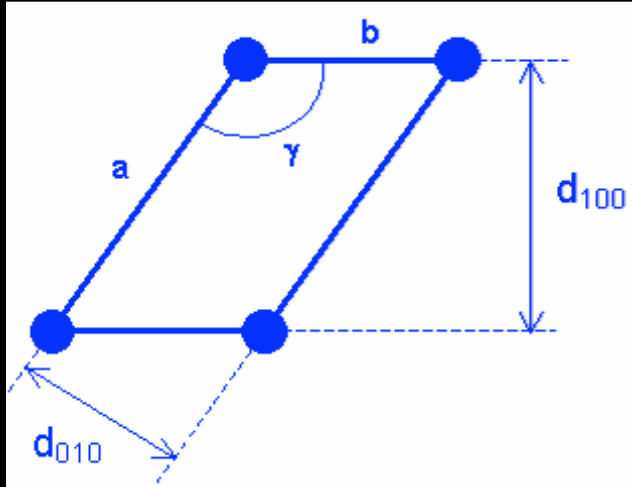
$$\left\{ \begin{array}{l} \mathbf{a} \perp \mathbf{b}^* \text{ and } \mathbf{c}^* \\ \mathbf{b} \perp \mathbf{c}^* \text{ and } \mathbf{a}^* \\ \mathbf{c} \perp \mathbf{a}^* \text{ and } \mathbf{b}^* \end{array} \right.$$

and

$$\left\{ \begin{array}{l} \mathbf{a}^* = |\mathbf{a}^*| = 1/d_{100} \\ \mathbf{b}^* = |\mathbf{b}^*| = 1/d_{010} \\ \mathbf{c}^* = |\mathbf{c}^*| = 1/d_{001} \end{array} \right.$$



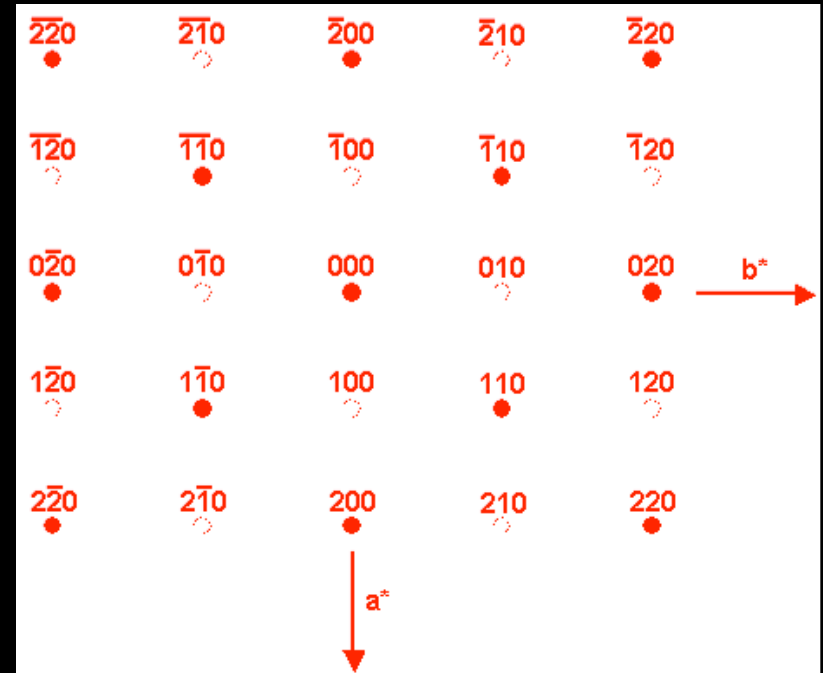
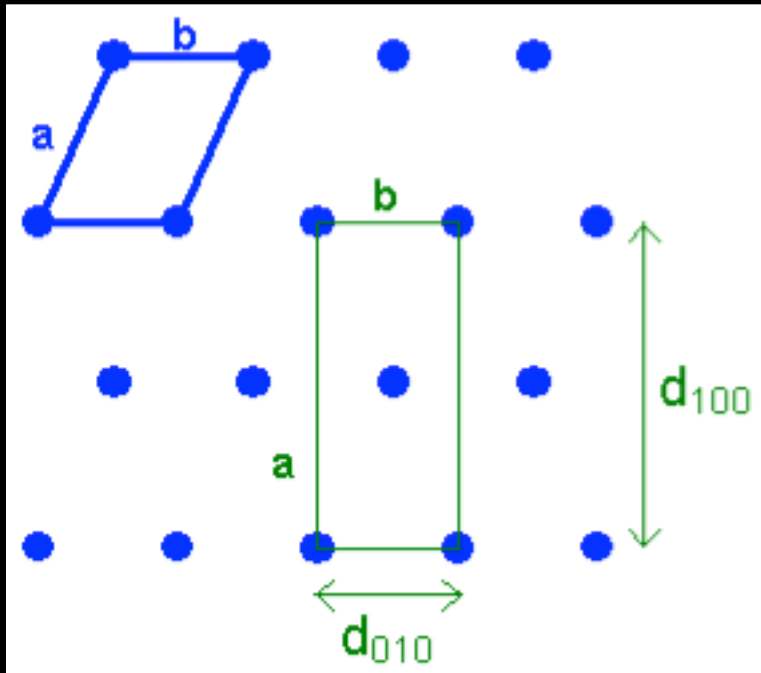
Primitive parallelogram lattice



N.B. In all crystal systems :

$$\begin{cases} \mathbf{a}^* \perp \mathbf{b} \text{ and } \mathbf{c} \\ \mathbf{b}^* \perp \mathbf{c} \text{ and } \mathbf{a} \\ \mathbf{c}^* \perp \mathbf{a} \text{ and } \mathbf{b} \end{cases} \quad \mathbf{and} \quad \begin{cases} \mathbf{a} \perp \mathbf{b}^* \text{ and } \mathbf{c}^* \\ \mathbf{b} \perp \mathbf{c}^* \text{ and } \mathbf{a}^* \\ \mathbf{c} \perp \mathbf{a}^* \text{ and } \mathbf{b}^* \end{cases} \quad \mathbf{and} \quad \begin{cases} a^* = |\mathbf{a}^*| = 1/d_{100} \\ b^* = |\mathbf{b}^*| = 1/d_{010} \\ c^* = |\mathbf{c}^*| = 1/d_{001} \end{cases}$$

Centered rectangular lattice



N.B. In all crystal systems :

$$\begin{cases} \mathbf{a}^* \perp \mathbf{b} \text{ and } \mathbf{c} \\ \mathbf{b}^* \perp \mathbf{c} \text{ and } \mathbf{a} \\ \mathbf{c}^* \perp \mathbf{a} \text{ and } \mathbf{b} \end{cases} \quad \mathbf{and} \quad \begin{cases} \mathbf{a} \perp \mathbf{b}^* \text{ and } \mathbf{c}^* \\ \mathbf{b} \perp \mathbf{c}^* \text{ and } \mathbf{a}^* \\ \mathbf{c} \perp \mathbf{a}^* \text{ and } \mathbf{b}^* \end{cases} \quad \mathbf{and} \quad \begin{cases} a^* = |\mathbf{a}^*| = 1/d_{100} \\ b^* = |\mathbf{b}^*| = 1/d_{010} \\ c^* = |\mathbf{c}^*| = 1/d_{001} \end{cases}$$

A family of lattice planes in real space
corresponds to
a point in reciprocal space.

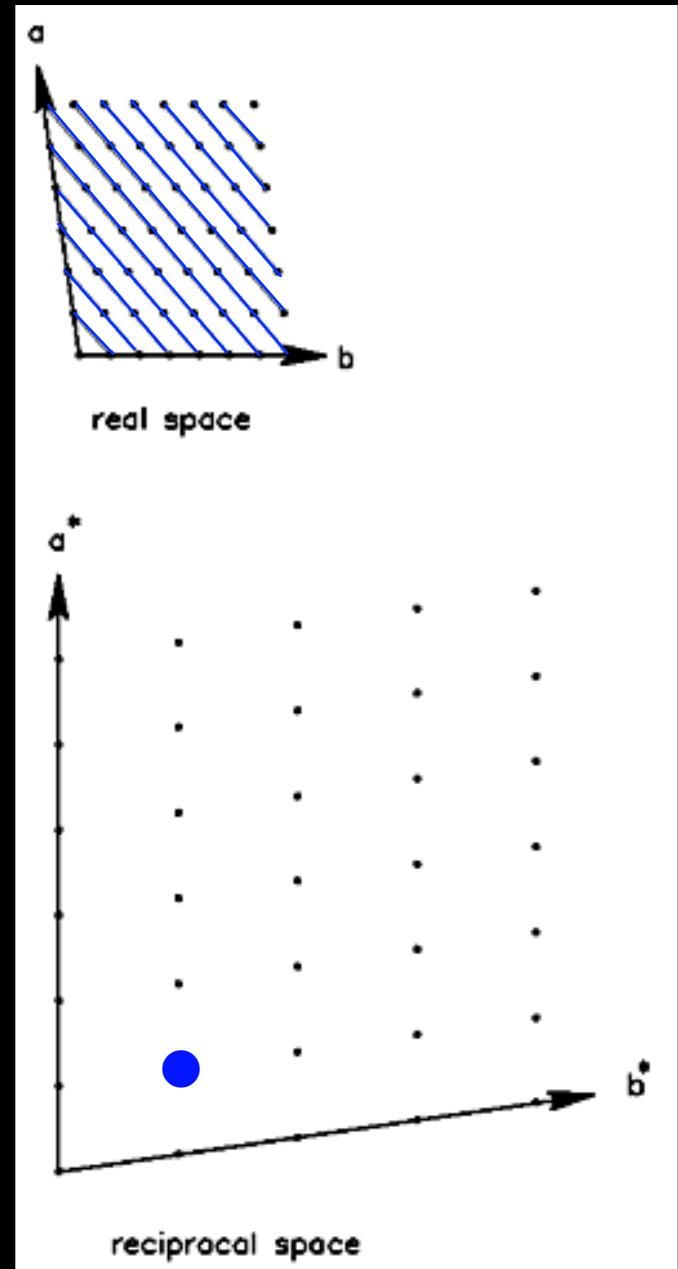
110 lattice planes

N.B. In all crystal systems:

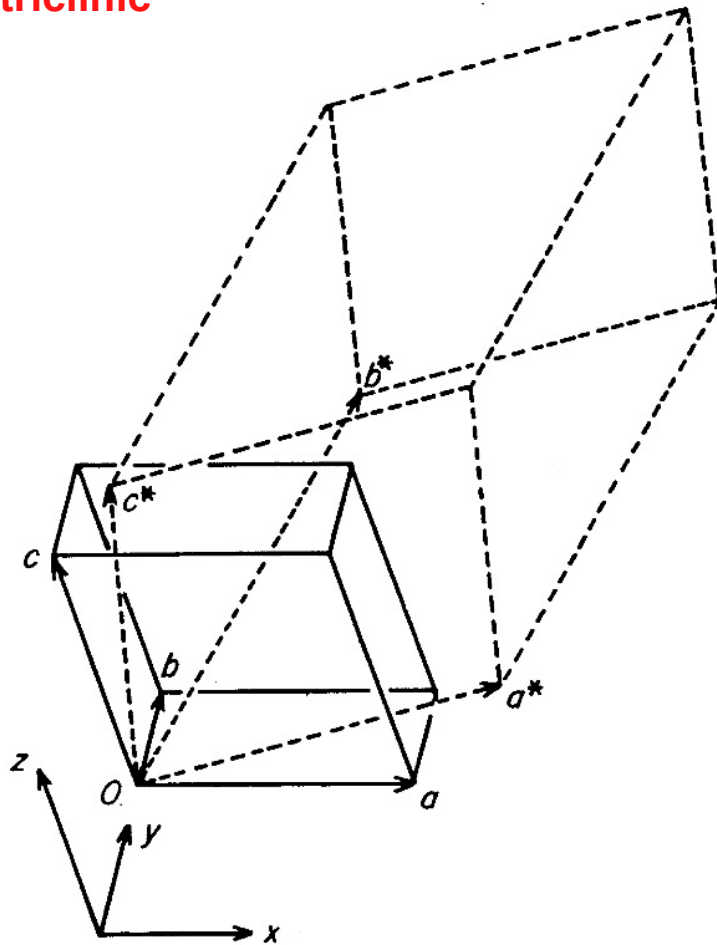
$$\left\{ \begin{array}{l} \mathbf{a}^* \perp \mathbf{b} \text{ and } \mathbf{c} \\ \mathbf{b}^* \perp \mathbf{c} \text{ and } \mathbf{a} \\ \mathbf{c}^* \perp \mathbf{a} \text{ and } \mathbf{b} \end{array} \right. \quad \text{and} \quad \left\{ \begin{array}{l} \mathbf{a} \perp \mathbf{b}^* \text{ and } \mathbf{c}^* \\ \mathbf{b} \perp \mathbf{c}^* \text{ and } \mathbf{a}^* \\ \mathbf{c} \perp \mathbf{a}^* \text{ and } \mathbf{b}^* \end{array} \right.$$

$$\text{and} \quad \left\{ \begin{array}{l} a^* = |\mathbf{a}^*| = 1/d_{100} \\ b^* = |\mathbf{b}^*| = 1/d_{010} \\ c^* = |\mathbf{c}^*| = 1/d_{001} \end{array} \right.$$

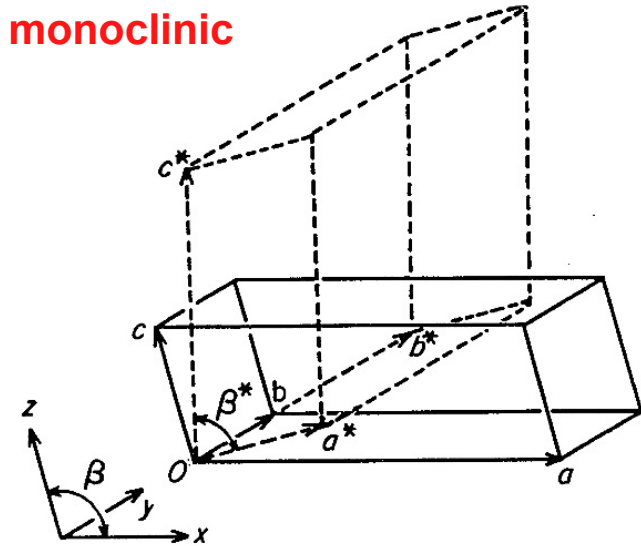
110 reciprocal lattice point



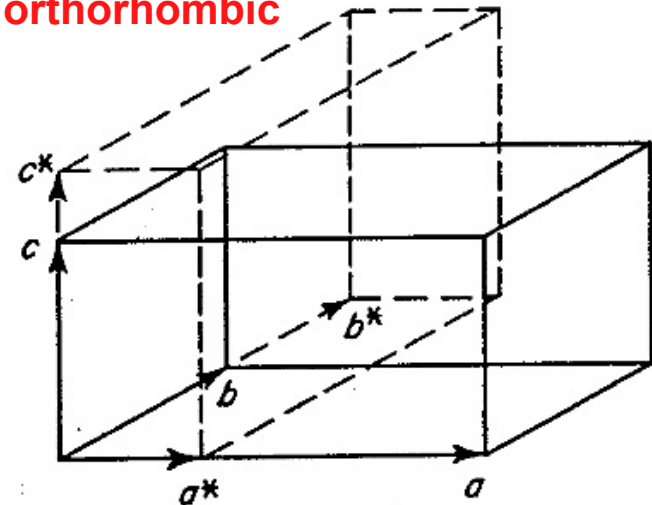
triclinic



monoclinic

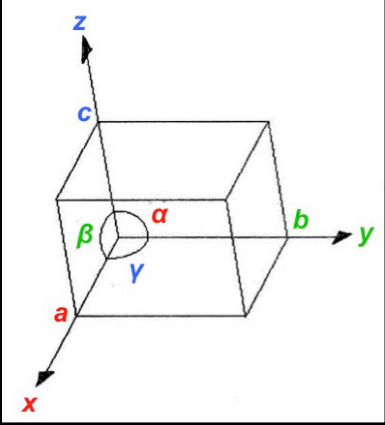


orthorhombic



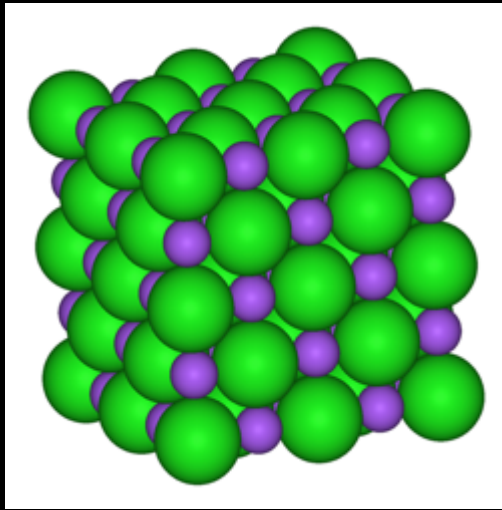
N.B. In all crystal systems :

$$\left\{ \begin{array}{l} \mathbf{a}^* \perp \mathbf{b} \text{ and } \mathbf{c} \\ \mathbf{b}^* \perp \mathbf{c} \text{ and } \mathbf{a} \\ \mathbf{c}^* \perp \mathbf{a} \text{ and } \mathbf{b} \end{array} \right. \quad \text{and} \quad \left\{ \begin{array}{l} \mathbf{a} \perp \mathbf{b}^* \text{ and } \mathbf{c}^* \\ \mathbf{b} \perp \mathbf{c}^* \text{ and } \mathbf{a}^* \\ \mathbf{c} \perp \mathbf{a}^* \text{ and } \mathbf{b}^* \end{array} \right. \quad \text{and} \quad \left\{ \begin{array}{l} a^* = |\mathbf{a}^*| = 1/d_{100} \\ b^* = |\mathbf{b}^*| = 1/d_{010} \\ c^* = |\mathbf{c}^*| = 1/d_{001} \end{array} \right.$$

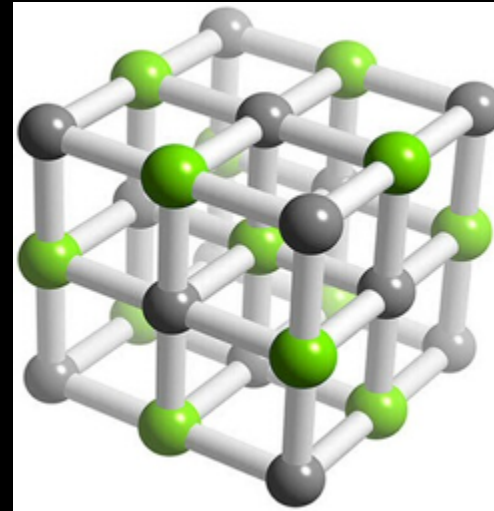


The Braggs' first crystal structures (1913)

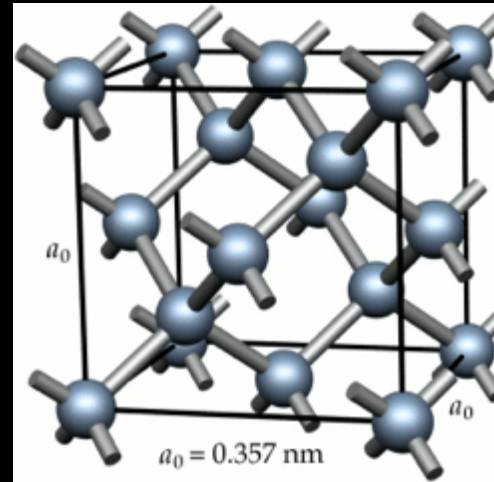
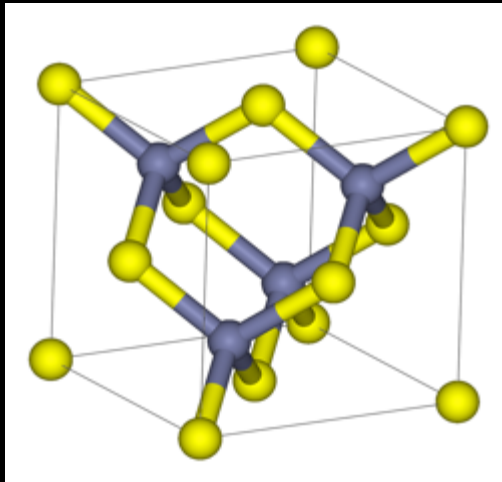
NaCl
rock salt
halite



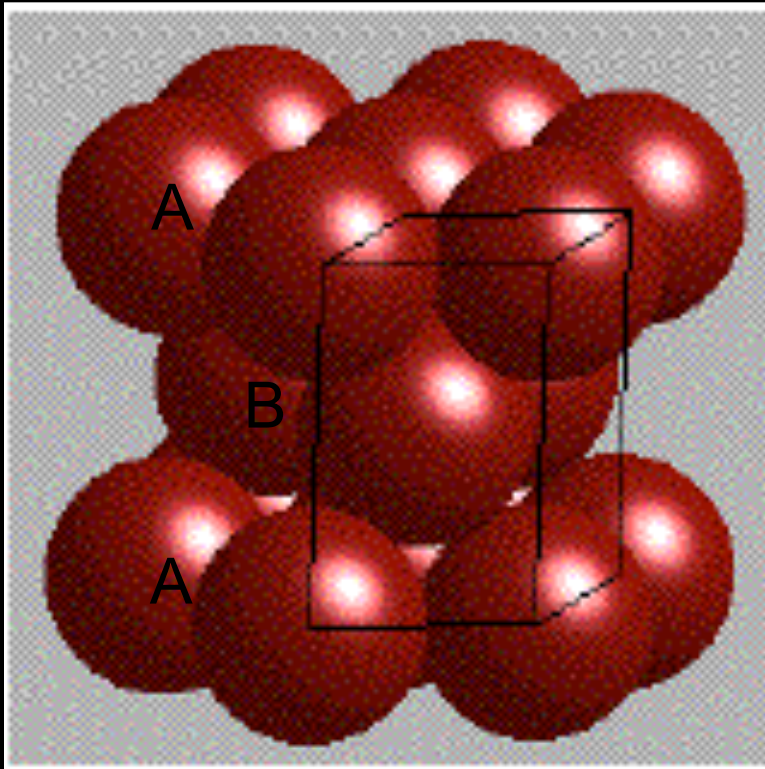
KCl
sylvite



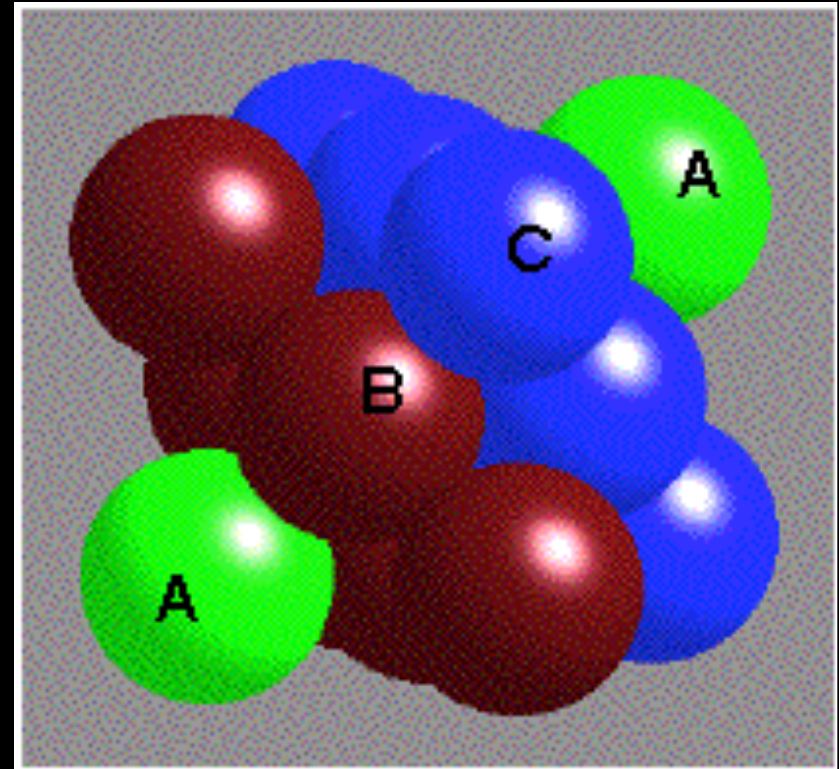
ZnS
zinc blende
spahlerite



Closest packing of spheres (Kepler's conjecture, 1611)

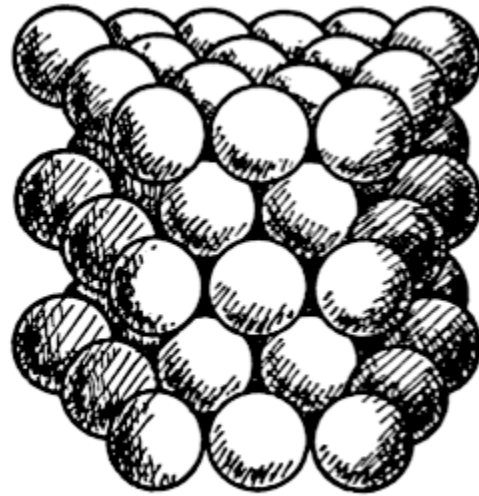


HCP
ABA...

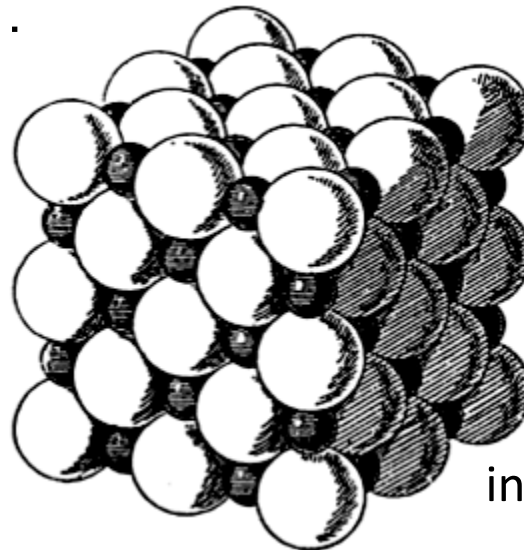
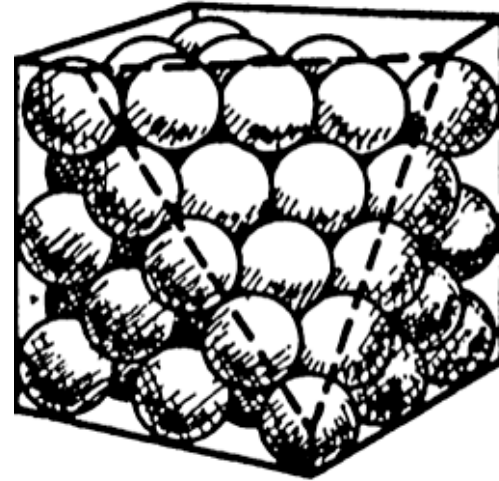


CCP = FCC
ABCA

William Barlow (1898) and William Jackson Pope (1906)

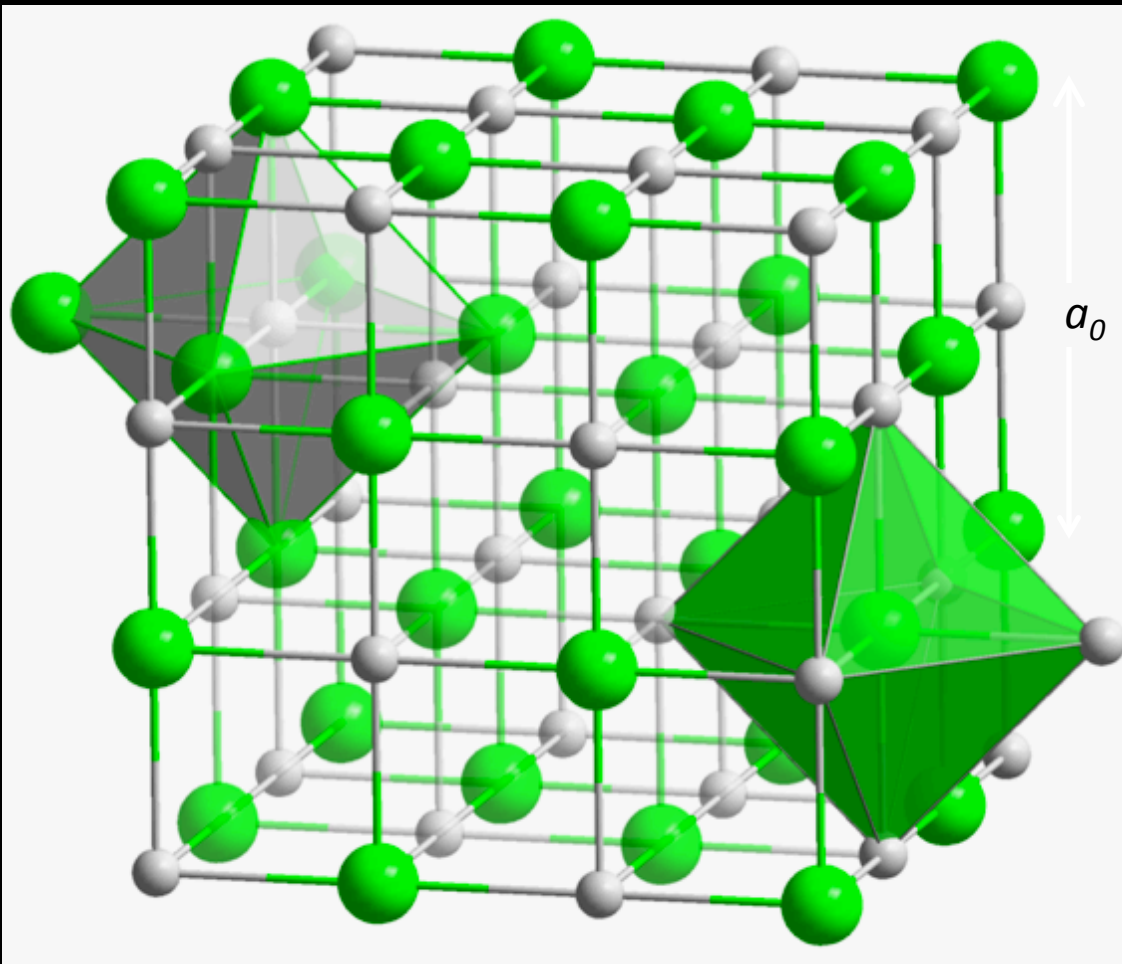


hcp
ABA...



NaCl
interpenetrating
fcc

Hypothetical crystal structures based on sphere-packing



$$\sin \theta_{hkl} = \frac{n\lambda}{2d_{hkl}}$$

$$d_{hkl} = a_0 / \sqrt{h^2 + k^2 + l^2}$$

$$d_{100} = a_0$$

$$d_{110} = a_0 / \sqrt{2} = 0.707 a_0$$

$$d_{111} = a_0 / \sqrt{3} = 0.577 a_0$$

$$d_{200} = a_0 / 2$$

$$d_{210} = a_0 / \sqrt{5} = 0.447 a_0$$

$$d_{211} = a_0 / \sqrt{6} = 0.408 a_0$$

⋮

$$a_0 = \sqrt[3]{V_{\text{cell}}} = \sqrt[3]{\frac{Z_{\text{cell}} M_r}{\rho_m N_A}}$$

$$= \sqrt[3]{\frac{4 \times 58.5 \times 10^{24} \text{ \AA}^3}{2.16 \times 6.02 \times 10^{23}}} = 5.63 \text{ \AA}$$

NaCl formula unit, $\text{Na}^+ (\text{Cl}^-)_{6 \times \frac{1}{6}}$, $\text{Cl}^- (\text{Na}^+)_{6 \times \frac{1}{6}}$, $M_r = 23.0 + 35.5 = 58.5 \text{ Da}$

Interpenetrating face-centered cubic unit cell, $Z_{\text{cell}} = 4\text{NaCl}$

$$\left(8 \times \frac{1}{8} \right) + \left(6 \times \frac{1}{2} \right) = 1 + 3 = 4 \left\{ \begin{array}{l} \text{Na}^+ \\ \text{Cl}^- \end{array} \right. \quad \text{and} \quad \left(12 \times \frac{1}{4} \right) + 1 = 3 + 1 = 4 \left\{ \begin{array}{l} \text{Cl}^- \\ \text{Na}^+ \end{array} \right.$$

corner Na^+
face Na^+
edge Cl^-
center Cl^-